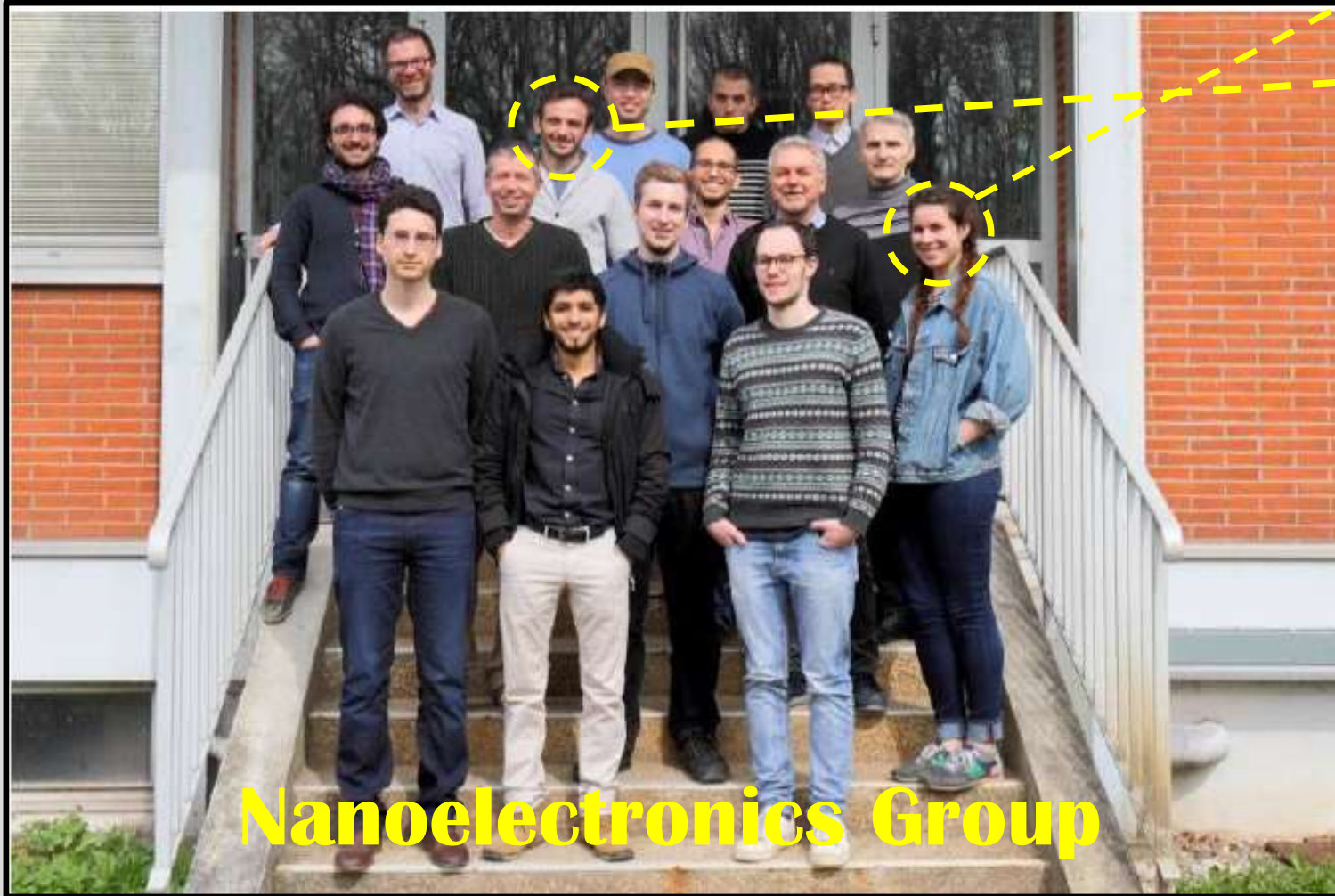


A Josephson relation for $e/3$ and $e/5$ fractionally charged with anyons



Nanoelectronics Group

Maelle Kapfer

Pređen Roulleau

D. C. G.

P. Jacques

@ NanoElectronics Group, CEA Saclay



D. Ritchie,

I. Farrer ,

@ Cambridge UK

OPEN POSITION
for 18-24 months
Post-doct.
(urgent)

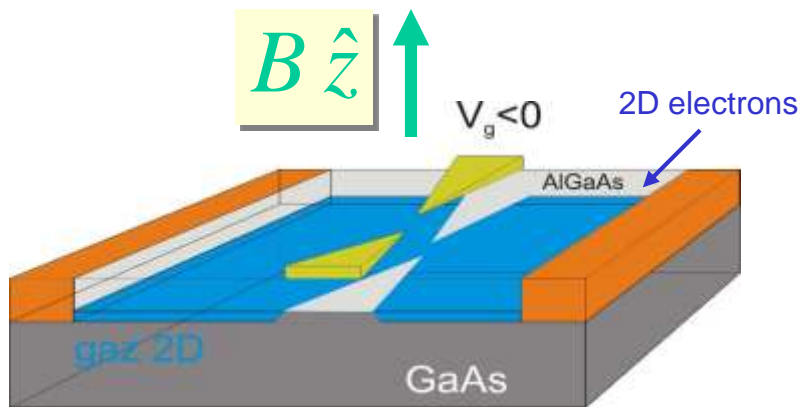
OUTLINE

- Quantum Hall edge states and Fractional Quantum Hall Effect
- PHOTON-ASSISTED TRANSPORT
 - Photon-assisted processes
 - A JOSEPHSON Relation for Photon Assisted Shot Noise (PASN)
- Experimental Results
 - $e^*=e/3$
 - $e^*=e/5$
- CONCLUSION and PERSPECTIVES

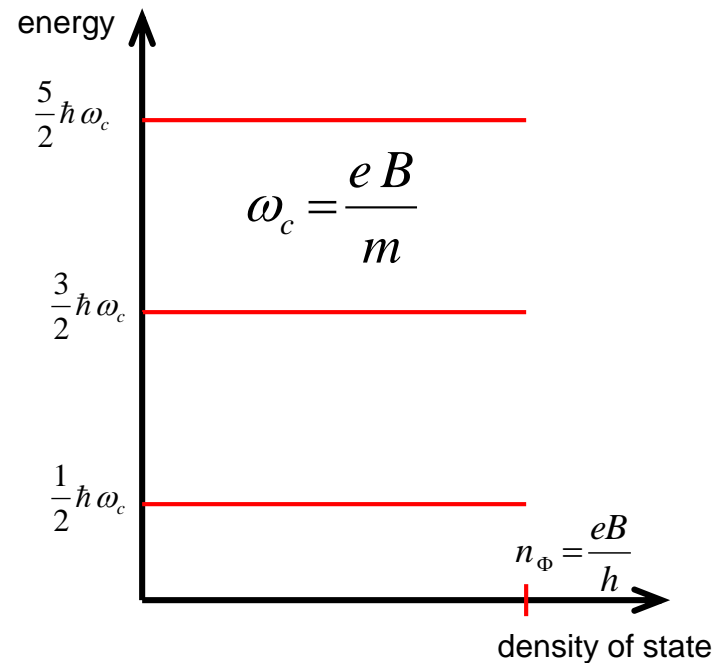
$$f_{J.} = \frac{e^*V}{h}$$

X. G. Wen (1991)

Quantum Hall Effect (QHE)



III-V semi-conductor heterojunction GaAs/GaAlAs



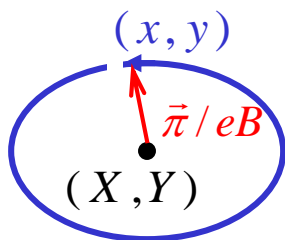
$$R_{Hall} = \frac{B}{en_s} = \frac{h}{e^2} \frac{1}{\nu = k}$$

$$H = \frac{1}{2m} (\vec{p} + e\vec{A})^2 = \frac{\vec{\pi}^2}{2m}$$

$$[\pi_x, \pi_y] = -i\hbar eB \quad \longrightarrow \quad E_n = \hbar\omega_c \left(n + \frac{1}{2} \right)$$

$$X = x - \frac{\pi_y}{eB}$$

$$Y = y + \frac{\pi_x}{eB}$$



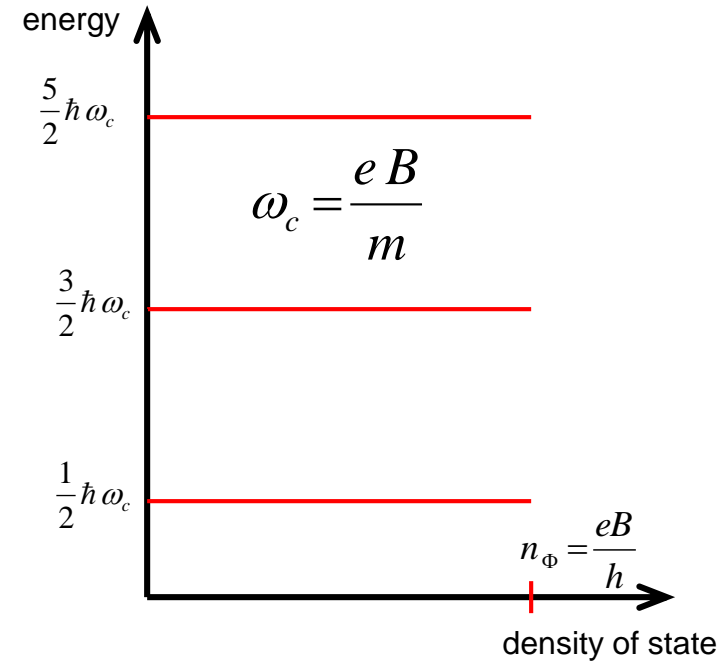
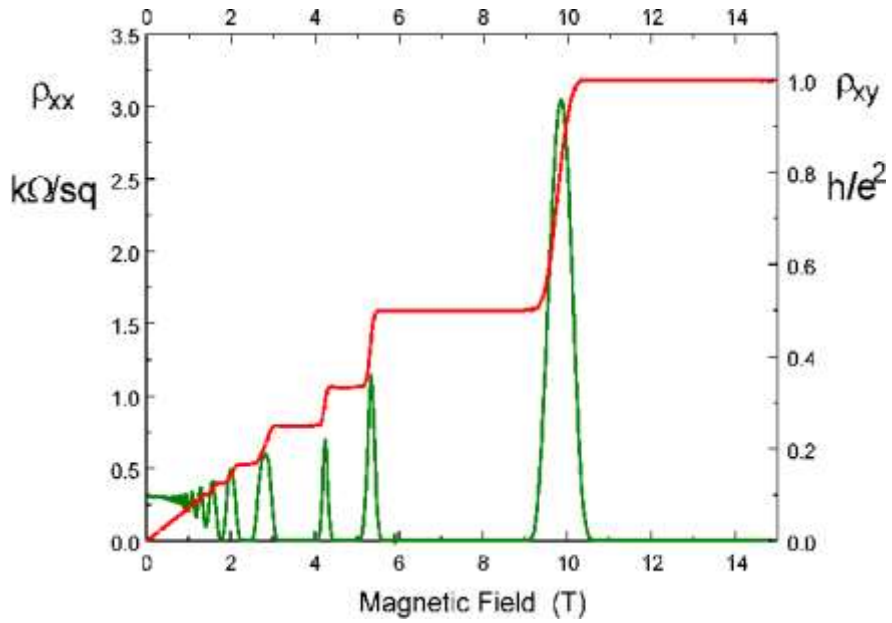
cyclotron motion

$$[X, Y] = -i\frac{\hbar}{eB} \quad \longrightarrow \quad B\Delta X.\Delta Y = \frac{h}{e}$$

cyclotron motion is frozen \rightarrow 1D dynamics

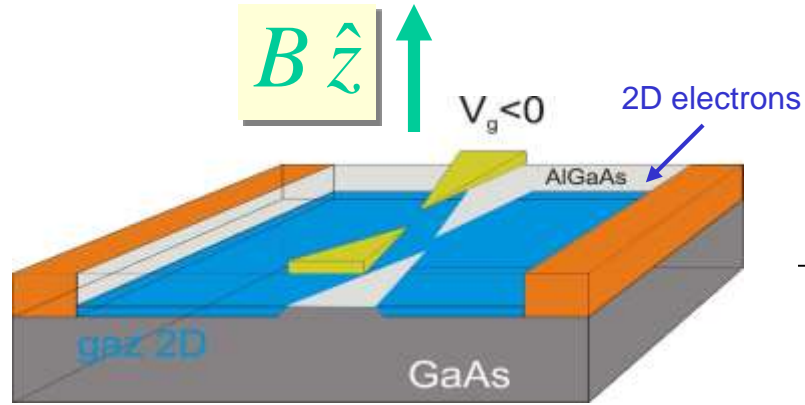
Integer Quantum Hall Effect (IQHE)

$$R_{\text{hall}} = (h/e^2) \nu \quad \nu = 1, 2, 3, \dots$$

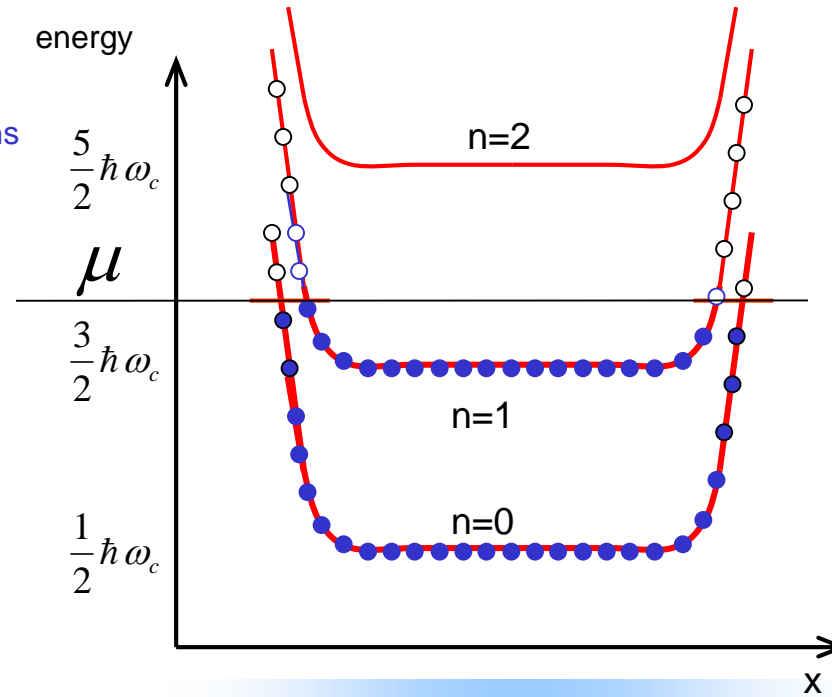


$$R_{\text{Hall}} = \frac{B}{e n_s} = \frac{h}{e^2} \frac{1}{\nu = k}$$

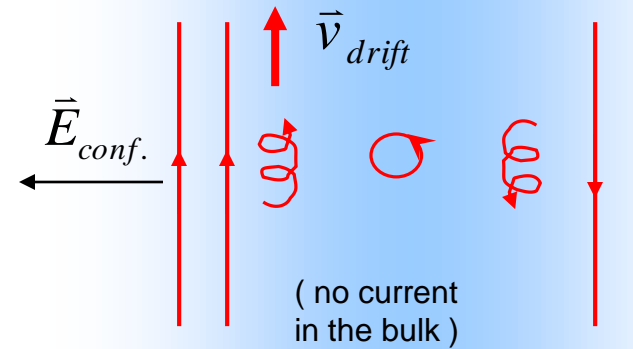
QHE and EDGE STATES



III-V semi-conductor heterojunction GaAs/GaAlAs



$$\vec{v}_{drift} = \frac{\vec{E}_{conf.}}{B} \times \hat{z}$$

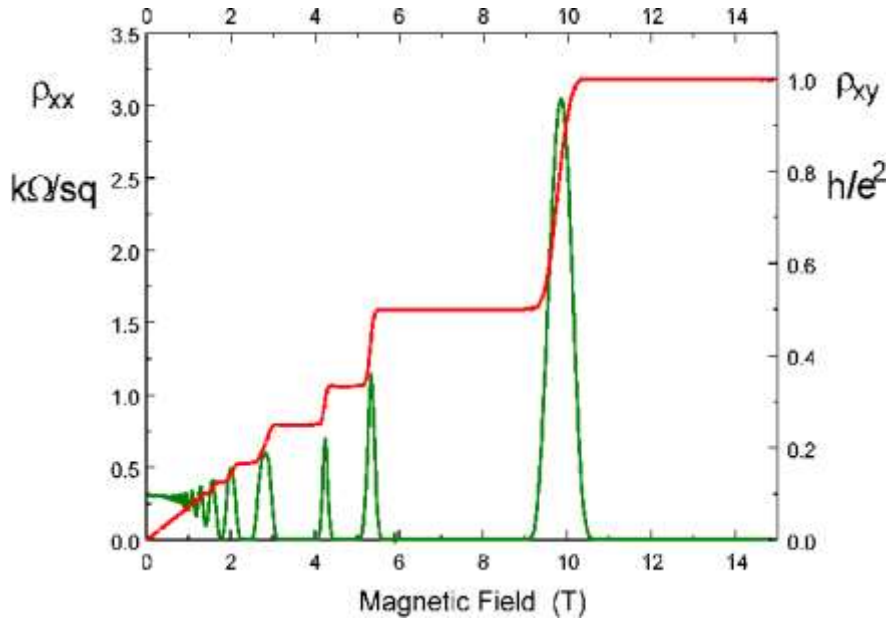


(edge current)

cyclotron motion drift → chiral 1D EDGE CHANNELS

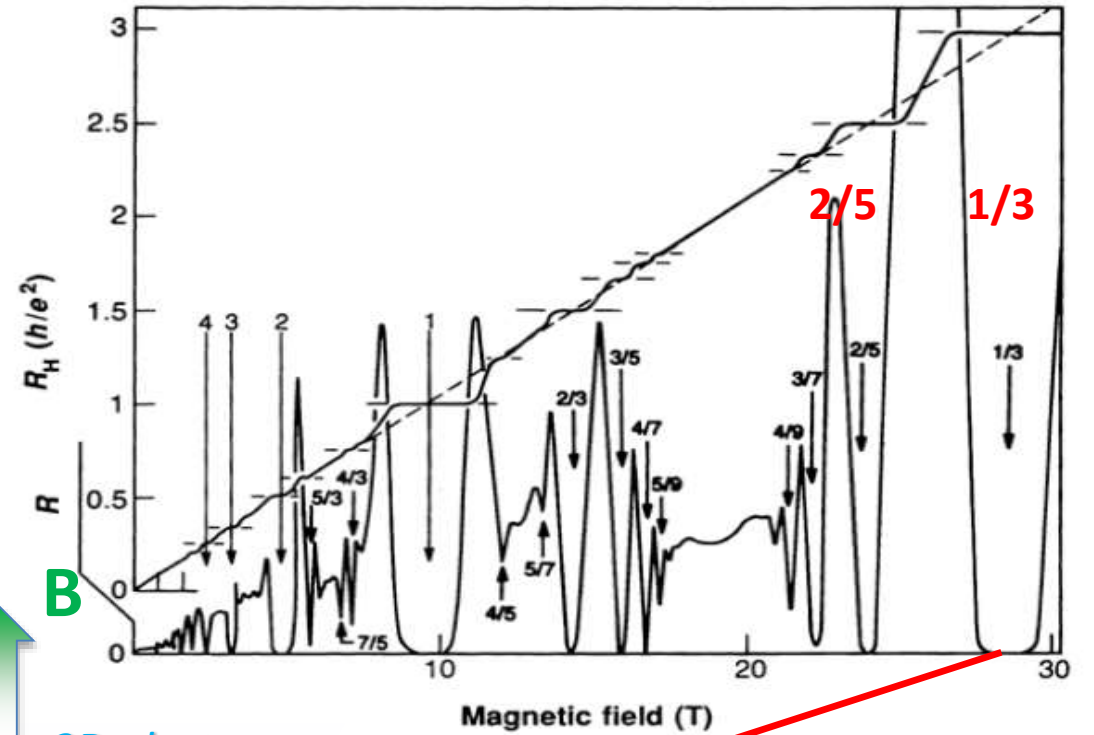
Integer Quantum Hall Effect (IQHE)

$$R_{\text{hall}} = (h/e^2)1/\nu \quad \nu = 1, 2, 3, \dots$$

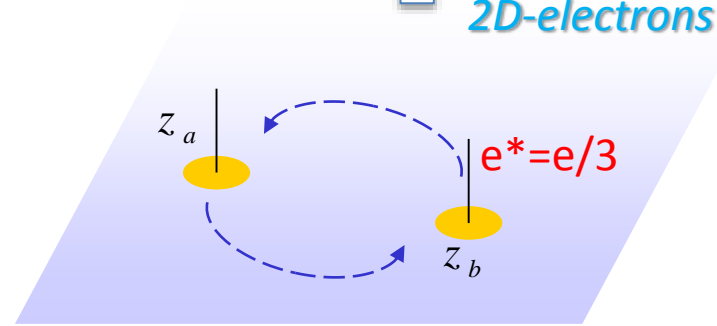


Fractional Quantum Hall Effect (FQHE)

$$R_{\text{hall}} = (h/e^2)1/\nu \quad \nu = 1/3, 2/5, 3/7, \dots 2/3, 3/5, 4/7, \dots$$



B
2D-electrons

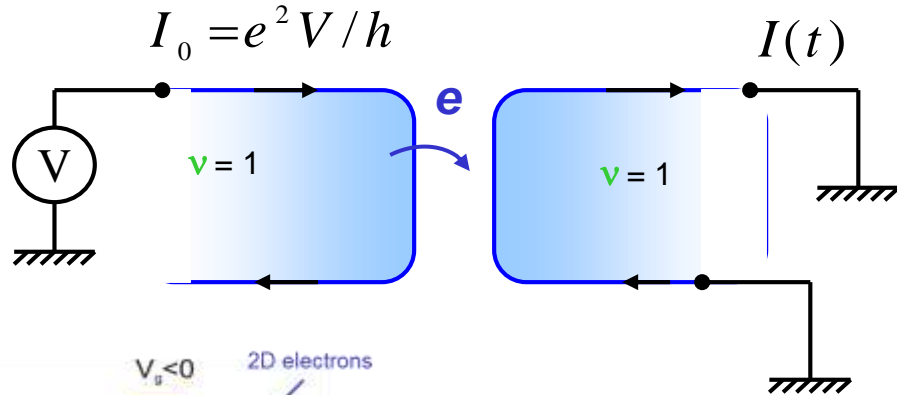


Anyons $\Psi(a,b) = e^{i\theta} \Psi(b,a) \quad \theta = 2\pi/3$

DC SHOT NOISE: Integer QHE

G. B. Lesovik,
JETP Letters 49, 594 (1989)

strong barrier :



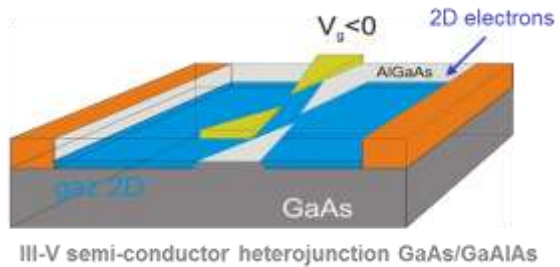
$$S_I = 2e I_0 D(1-D)$$

$$I_0 = e^2 V / h$$

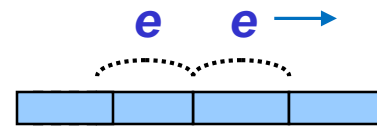
$$I_0 = I + I_B$$

transmitted (D)

reflected ($1-D$)



h/eV



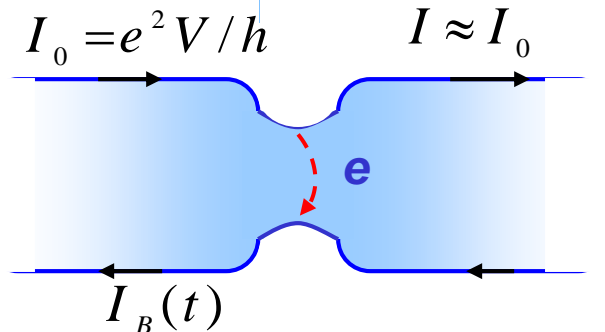
$$S_I = 2eI \quad D \ll 1$$

Schottky (1918)

(rarely transmitted electrons)

Poisson's statistics

weak barrier :

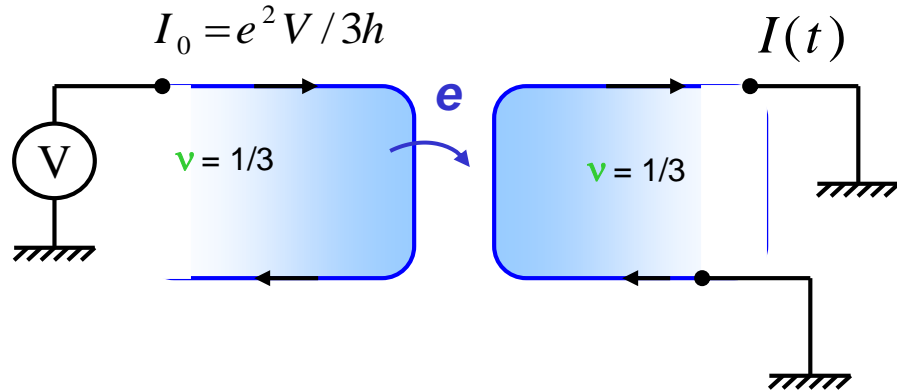


$$S_I = 2eI_B \quad D \approx 1$$

(rarely transmitted holes)

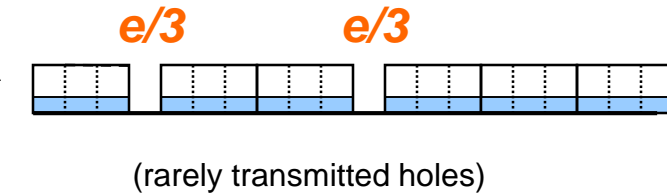
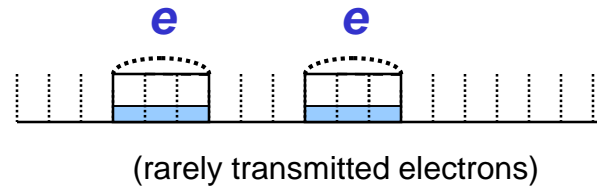
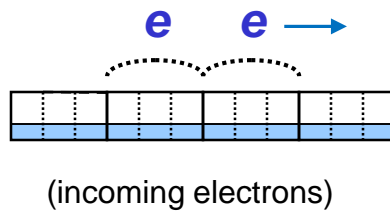
DC SHOT NOISE: FQHE

strong barrier :



$\nu=1/3$ Laughlin state

$$h/3eV$$



$$I_0 = e^2 V / 3h$$

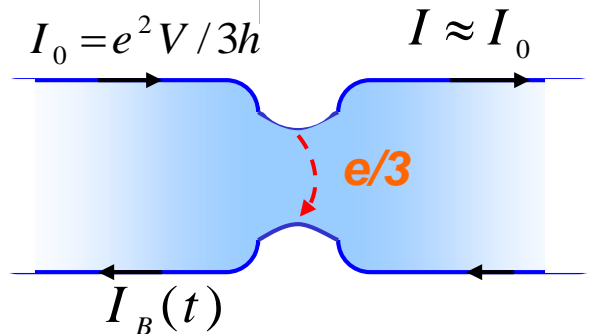
$$I_0 = I + I_B$$

transmitted (D)

reflected ($1-D$)

$$S_I = 2eI \quad D \ll 1$$

weak barrier :



$$S_I = 2 \frac{e}{3} I_B \quad D \approx 1$$

First observation:
CEA Saclay 1997
Weizmann 1997

derived from chiral-Luttinger liquid approach
(X.G. Wen 1995, C. Kane + M. Fisher 1994; Fendley, Ludwig + Saleur (1995))

Tunneling through a $\nu=2/5$ Jain FQHE state

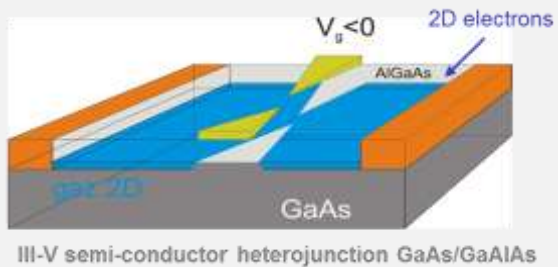
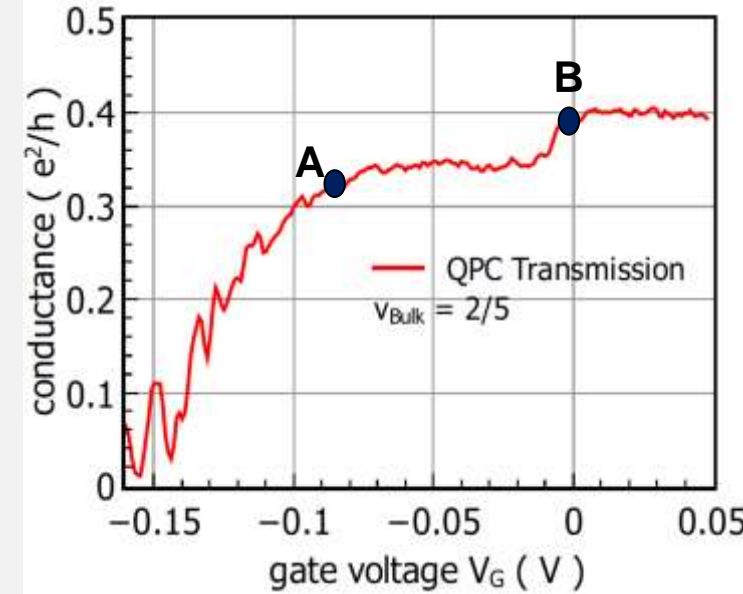
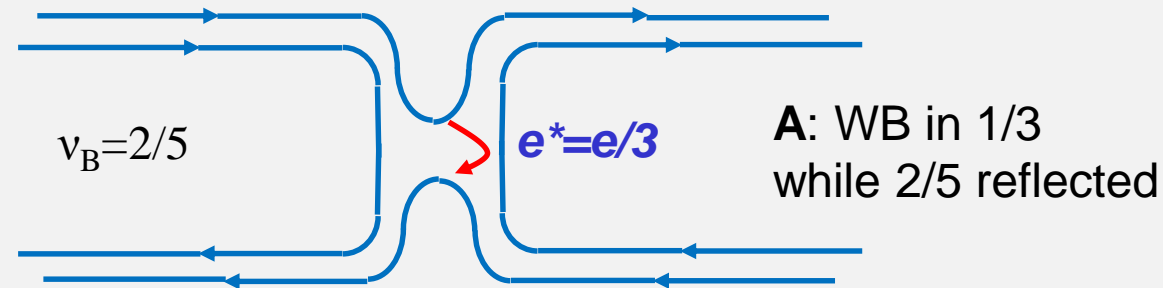
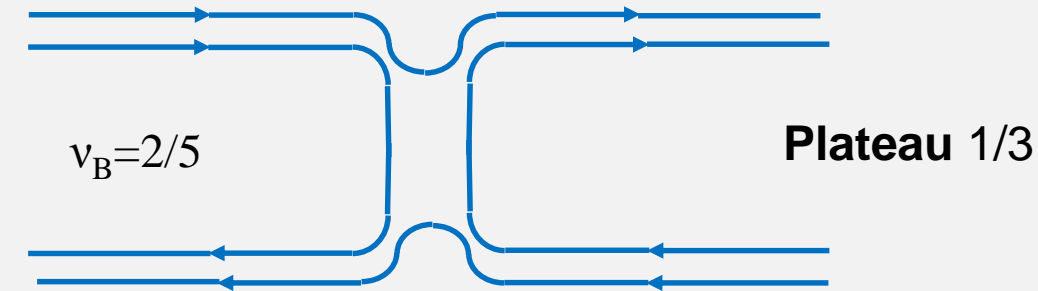
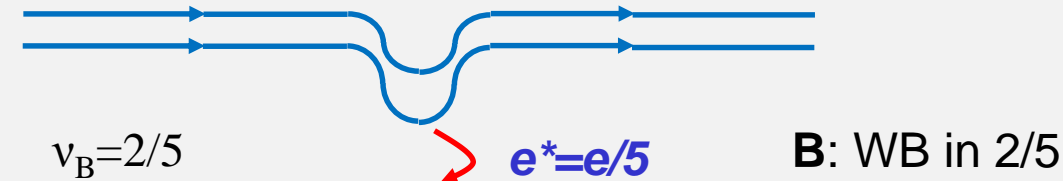
FQHE \rightarrow C-F. IQHE

$\nu = 1/3 \rightarrow \nu=1$

$\nu = 2/5 \rightarrow \nu=2$

$\nu = 3/7 \rightarrow \nu=3$

.... \rightarrow



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$$f_{J.} = \frac{e^*V}{h}$$

X. G. Wen (1991)

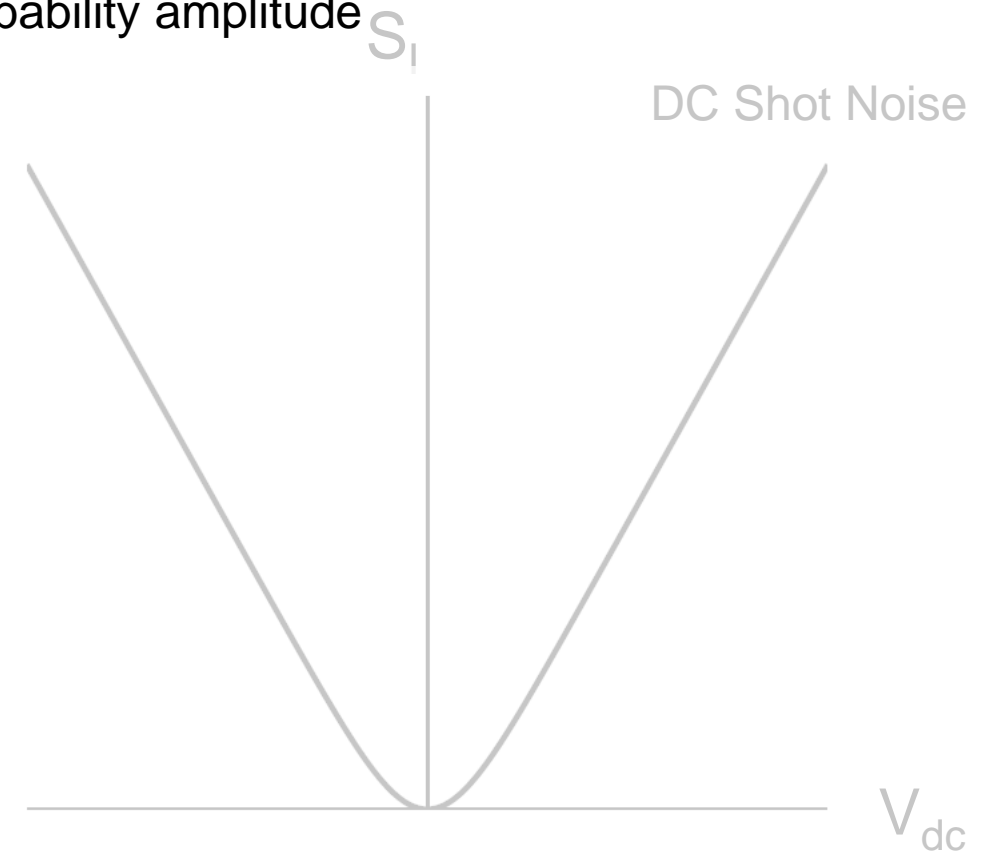
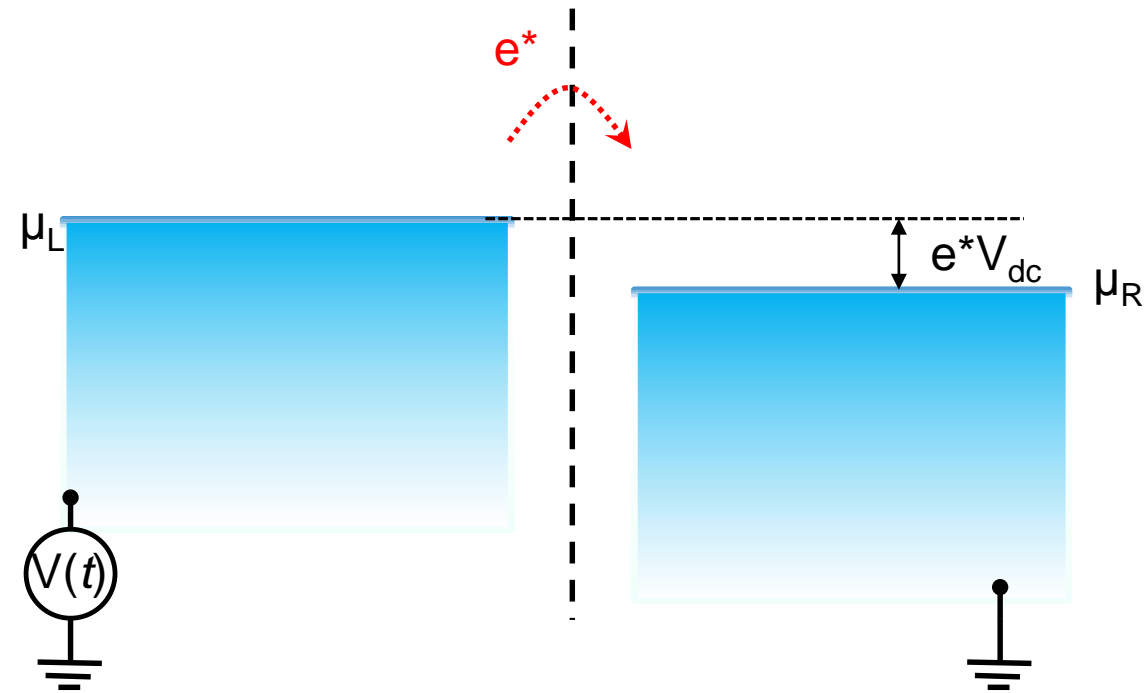
Photon-Assisted transport (weak coupling)

$$V(t) = V_{dc} + V_{ac} \cos(2\pi f t)$$

$$H_L \rightarrow H_L + e^* V_{ac} \cos(2\pi f t)$$

→ all carriers get extra time dependent phase: $\phi(t) = \frac{1}{\hbar} \int_{-\infty}^t e^* V_{ac}(t') dt'$

with : $\exp(-i\phi(t)) = \sum_l p_l e^{-i2\pi l f t}$ p_l : photo-absorption (Floquet) probability amplitude



Photon-Assisted Shot Noise (PASN)

$$V(t) = V_{dc} + V_{ac} \cos(2\pi f t)$$

$$H_L \rightarrow H_L + e^* V_{ac} \cos(2\pi f t)$$

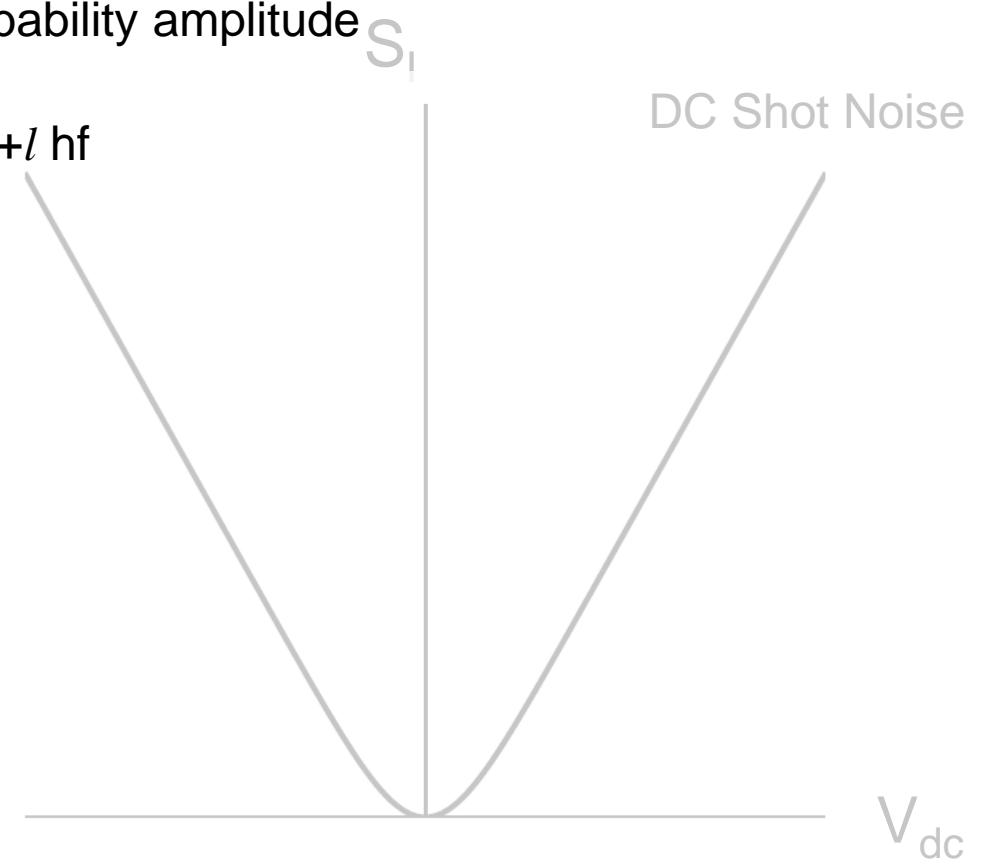
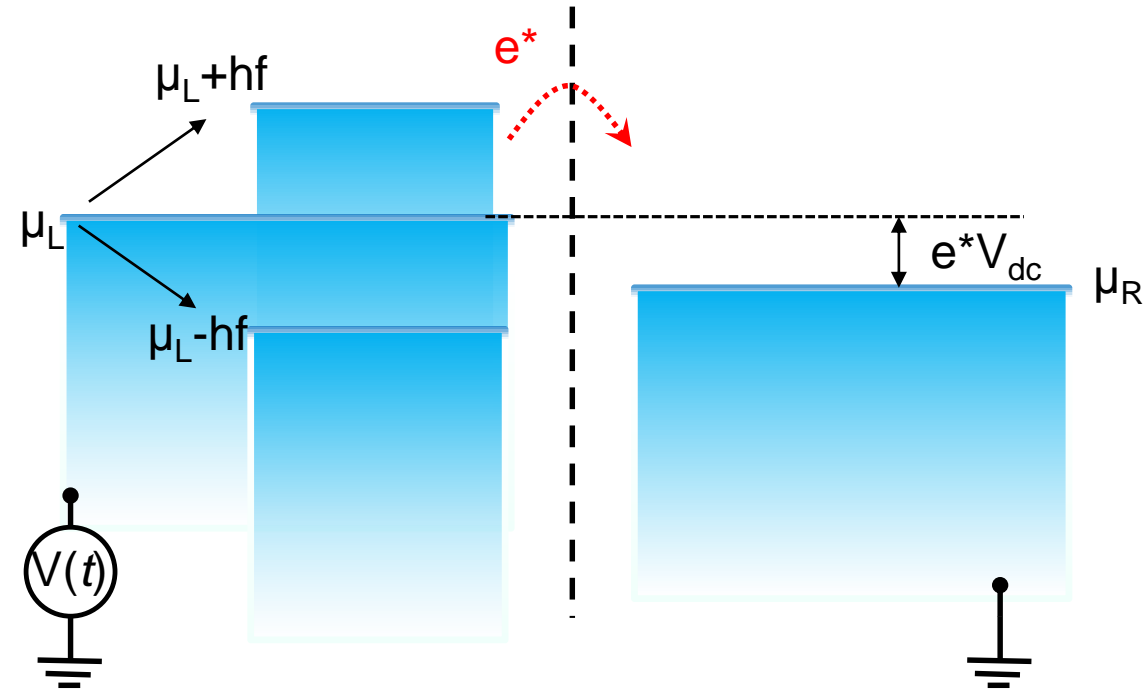
$$S_I^{PASN} = |p_0|^2 S_I^{DC}(V_{dc}) + |p_1|^2 S_I^{DC}(V_{dc} + hf/e^*) + |p_{-1}|^2 S_I^{DC}(V_{dc} - hf/e^*) + \dots$$

($e^*=e$) Lesovik and Levitov (1994)
 ($e^*=e/m$) Chamon and Wen (1995)

→ all carriers get extra time dependent phase: $\phi(t) = \frac{1}{\hbar} \int_{-\infty}^t e^* V_{ac}(t') dt'$

with : $\exp(-i\phi(t)) = \sum_l p_l e^{-i2\pi l f t}$ p_l : photo-absorption (Floquet) probability amplitude

global **energy scattering** for all left carrier energies ε shifted by $\varepsilon \rightarrow \varepsilon + l hf$



Photon-Assisted Shot Noise (PASN)

$$V(t) = V_{dc} + V_{ac} \cos(2\pi f t)$$

$$H_L \rightarrow H_L + e^* V_{ac} \cos(2\pi f t)$$

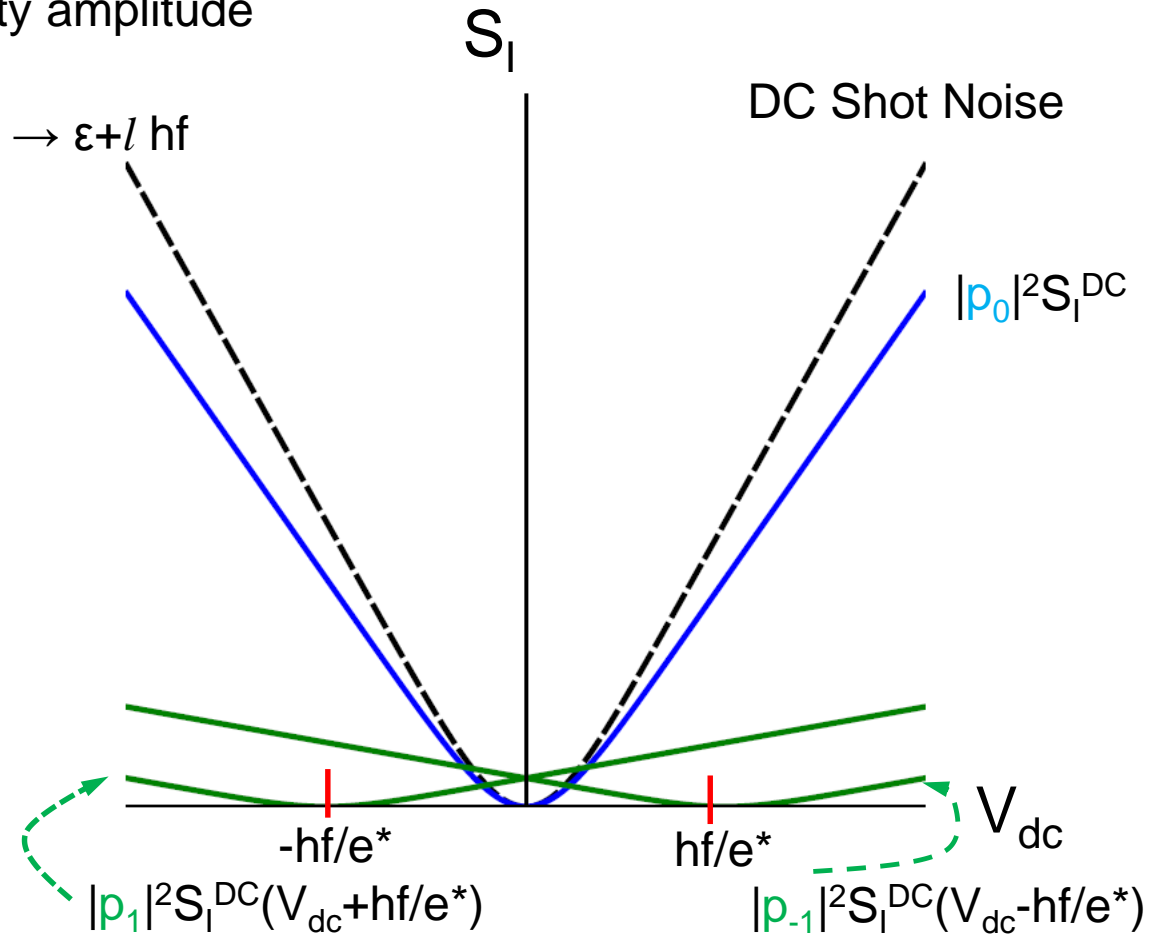
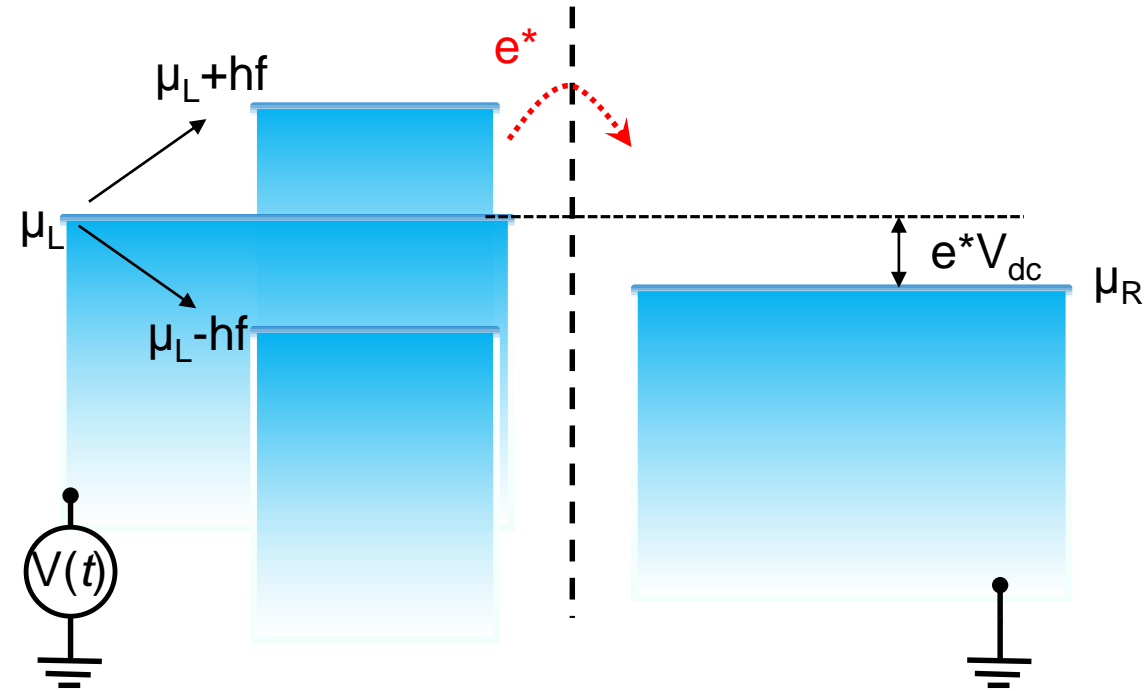
$$S_I^{PASN} = |p_0|^2 S_I^{DC}(V_{dc}) + |p_1|^2 S_I^{DC}(V_{dc} + hf/e^*) + |p_{-1}|^2 S_I^{DC}(V_{dc} - hf/e^*) + \dots$$

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global **energy scattering** for all left carrier energies ε shifted by $\varepsilon \rightarrow \varepsilon + l hf$



Photon-Assisted Shot Noise (PASN)

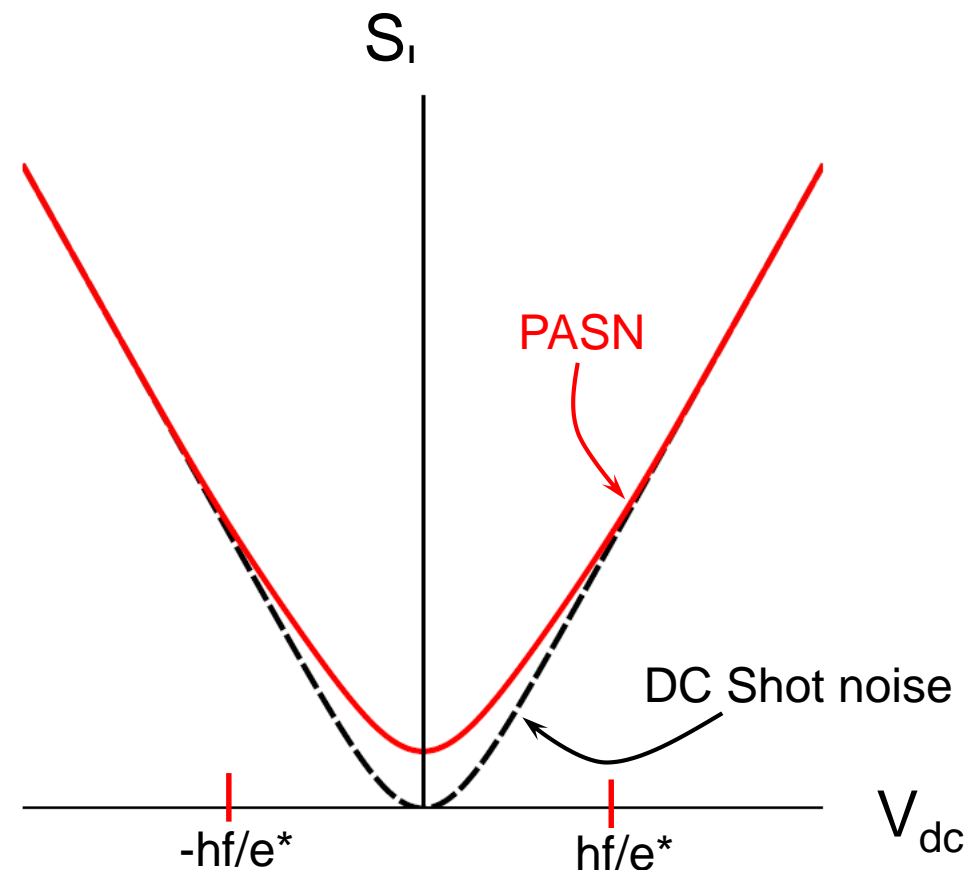
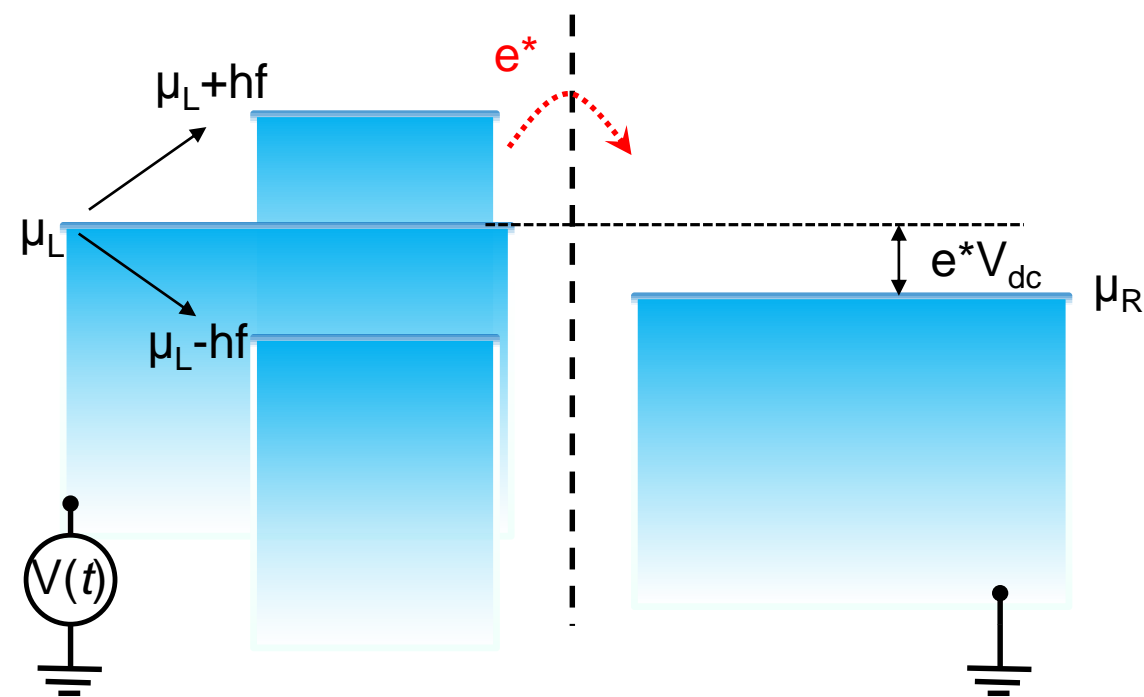
$$V(t) = V_{dc} + V_{ac} \cos(2\pi f t)$$

p_j : photo-absorption probability amplitude

μ_L shifted by $\mu_L \rightarrow \mu_L + l hf$ with probability $|p_l|^2$

$$|p_0|^2 + |p_1|^2 + |p_{-1}|^2 + \dots = 1$$

$$S_I^{PASN} = |p_0|^2 S_I^{DC}(V_{dc}) + |p_1|^2 S_I^{DC}(V_{dc} + hf/e^*) + |p_{-1}|^2 S_I^{DC}(V_{dc} - hf/e^*) + \dots$$



OUTLINE

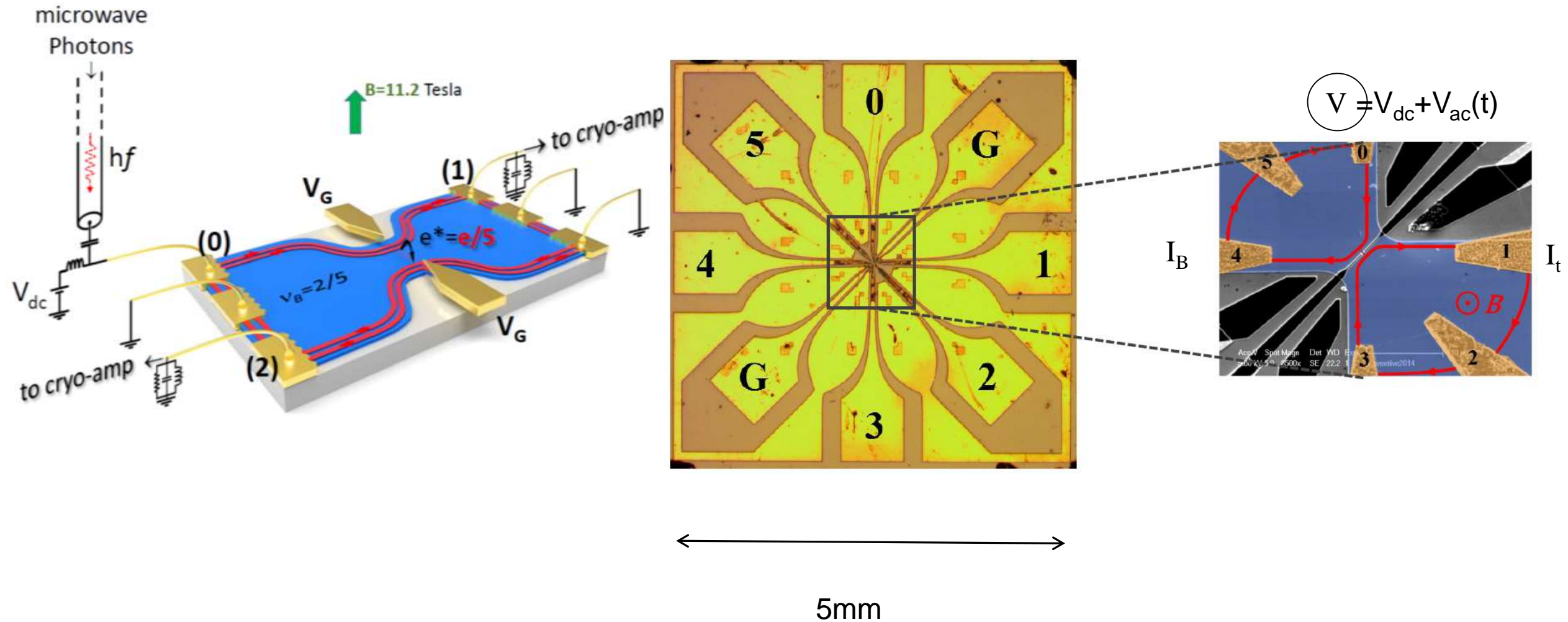
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$$f_{J.} = \frac{e^*V}{h}$$

X. G. Wen (1991)

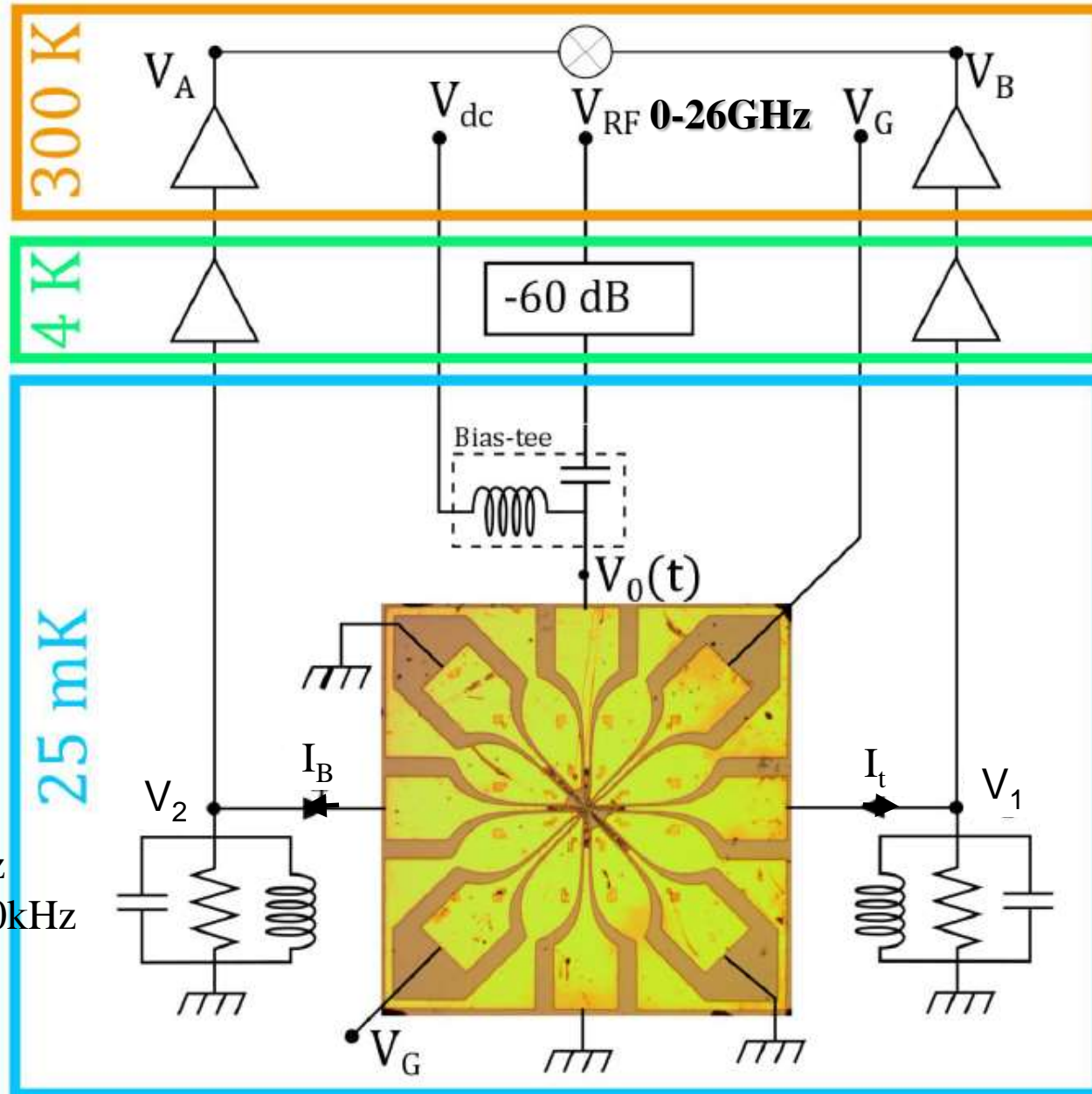
Experimental Set-up and samples

Samples: $n_s = 1.07 \cdot 10^{11} \text{ cm}^{-2}$ $\mu = 3 \cdot 10^6 \text{ cm}^2\text{V}^{-1}\text{s}^{-1}$ (from I. Farrer, D. Ritchie, Cambridge UK)



Experimental Set-up and samples

CROSS-SPECTRUM



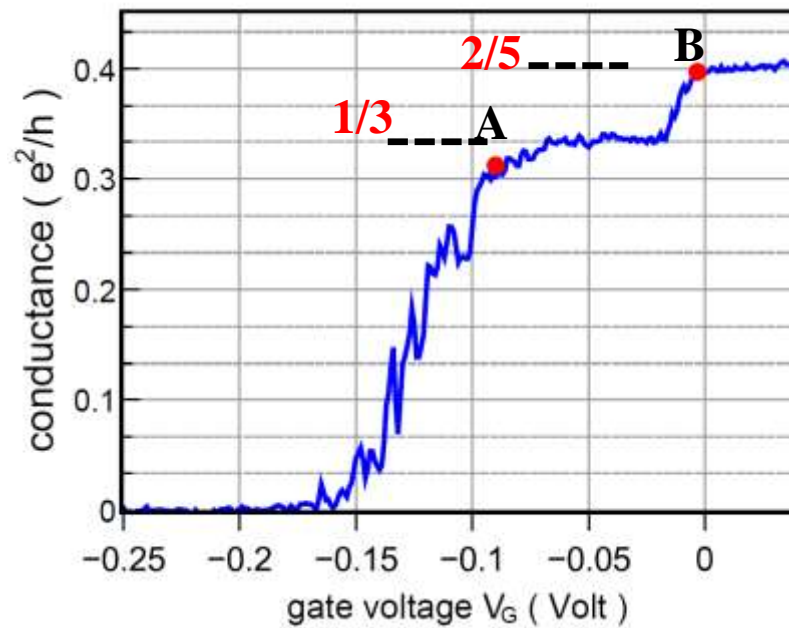
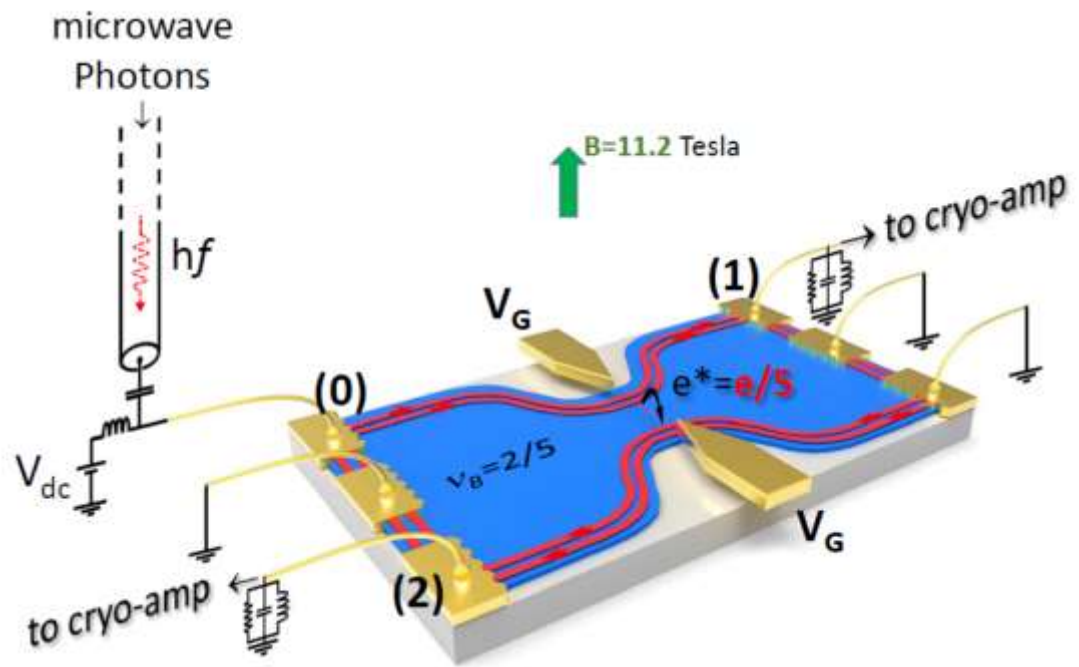
Helium-free Cryoconcept® cryostat



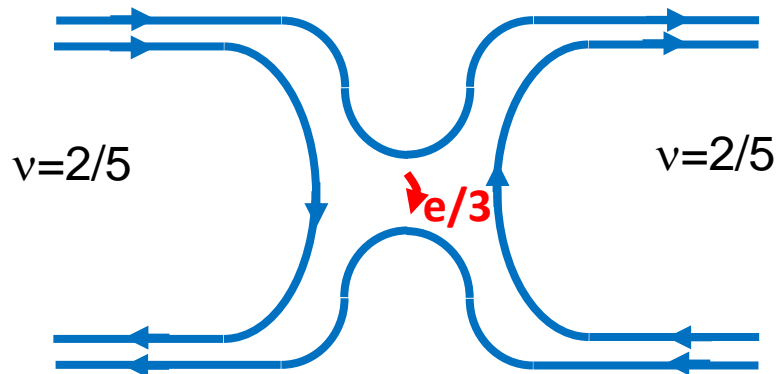
14 Tesla Dry Magnet
13mK base temperature

Home-made
Cryo-amp.
 $(0.22 \text{ nV})^2/\text{Hz}$

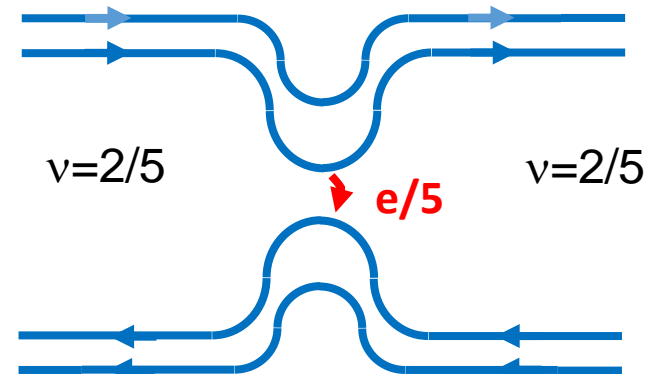
2.2 MHz
 $\Delta f \sim 150 \text{ kHz}$



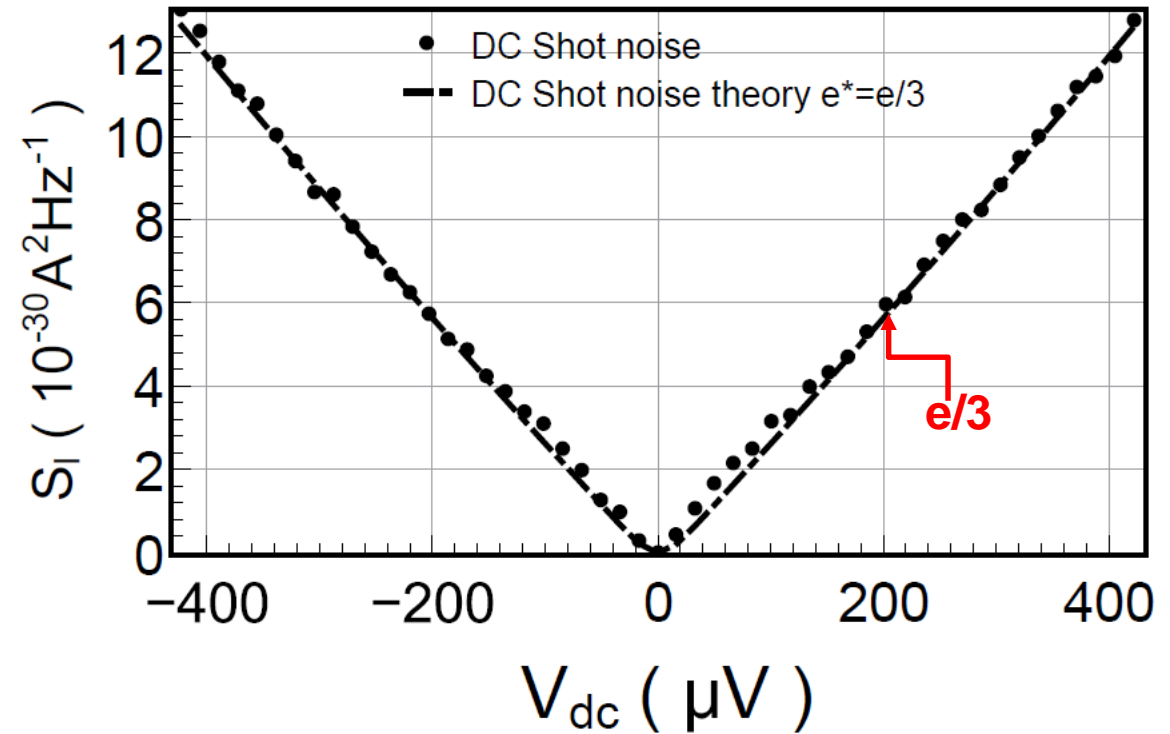
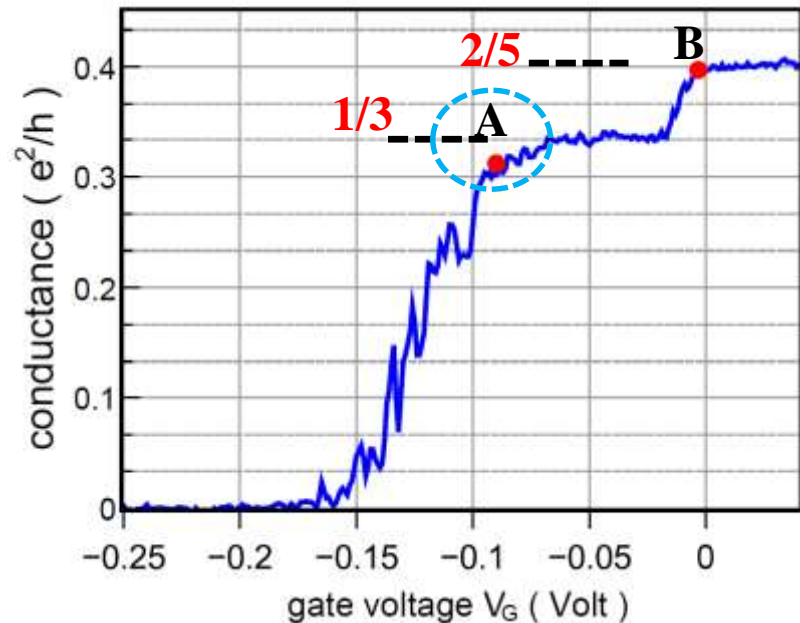
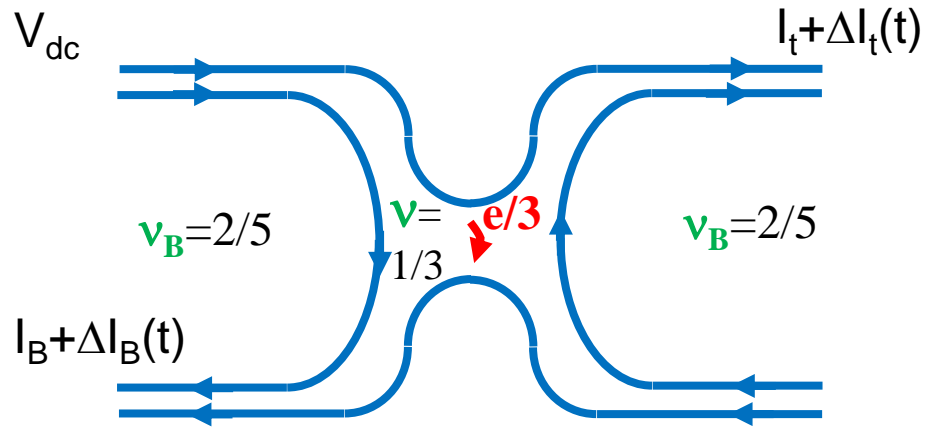
case A:



case B:



DC Shot noise for the 1/3-FQHE state



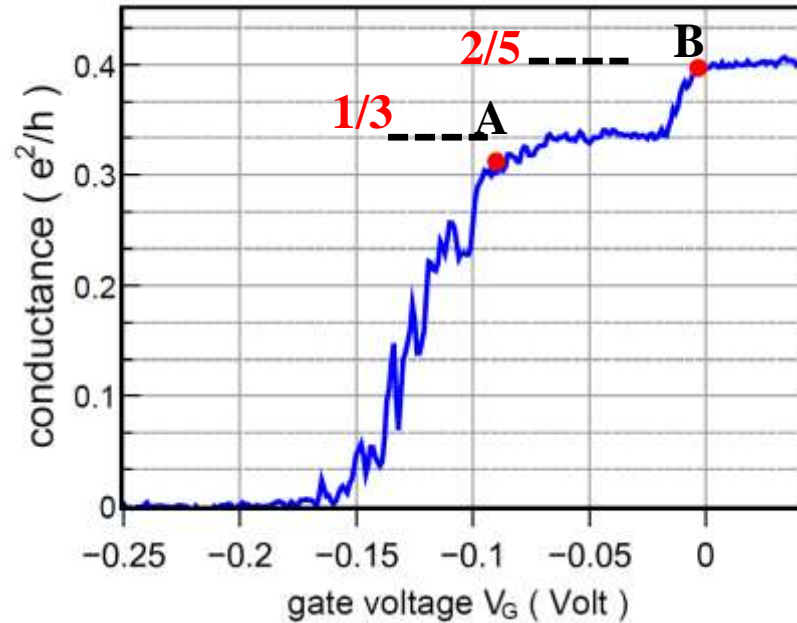
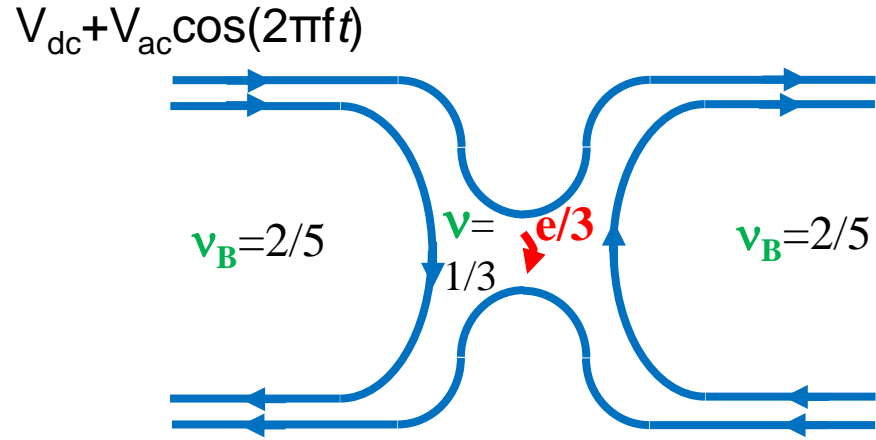
$$S_I^{DC} = 2e^* I_B \left[\coth \left(\frac{e^* V_{dc}}{2k_B T} \right) - \frac{2k_B T}{e^* V_{dc}} \right]$$

$e^* = e/3$!

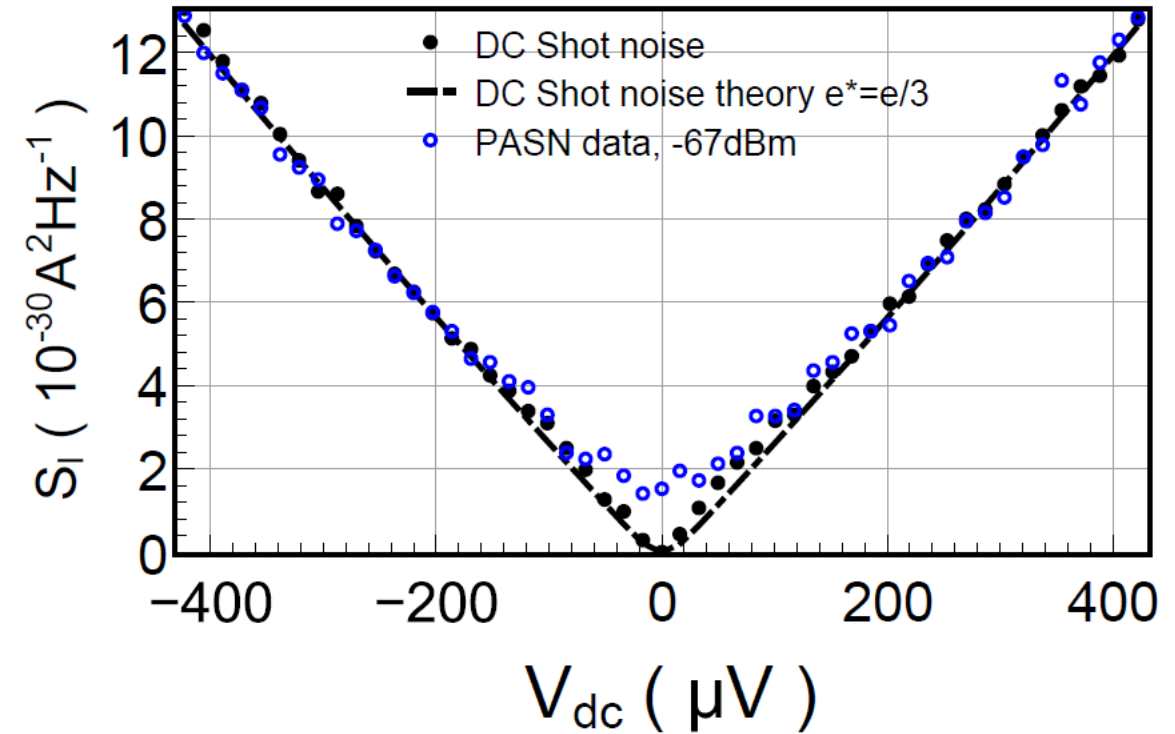
confirms '97-'98 experiments

(Saclay PRL 97, Weizmann Nat. 97 and 99, NTT 2015)

Photon-Assisted Shot Noise for the 1/3-FQHE state



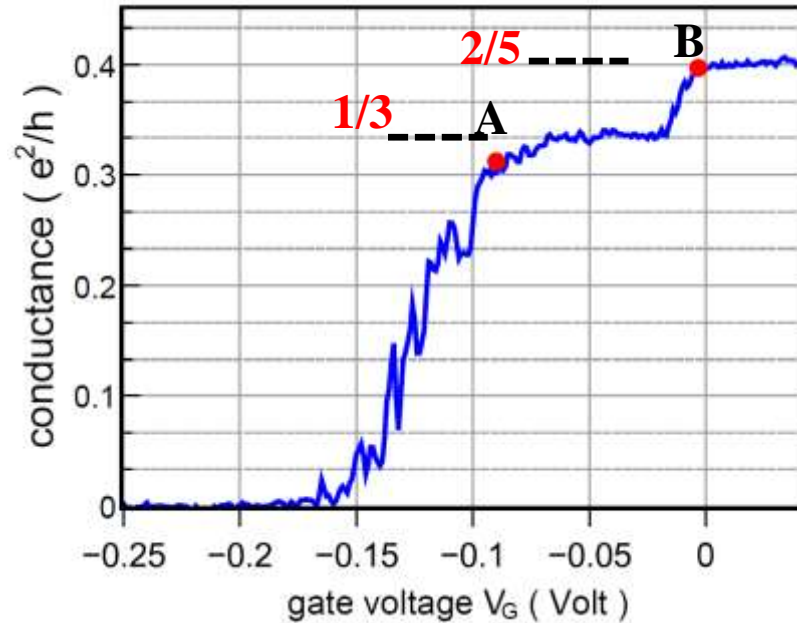
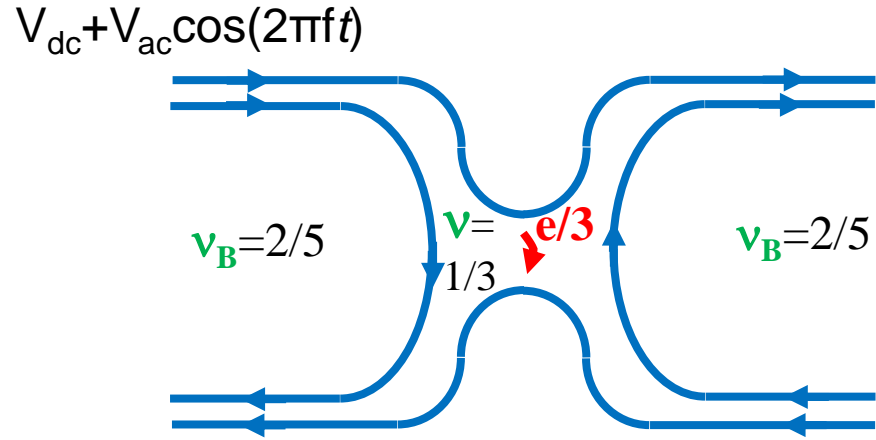
$f = 22 \text{ GHz}$



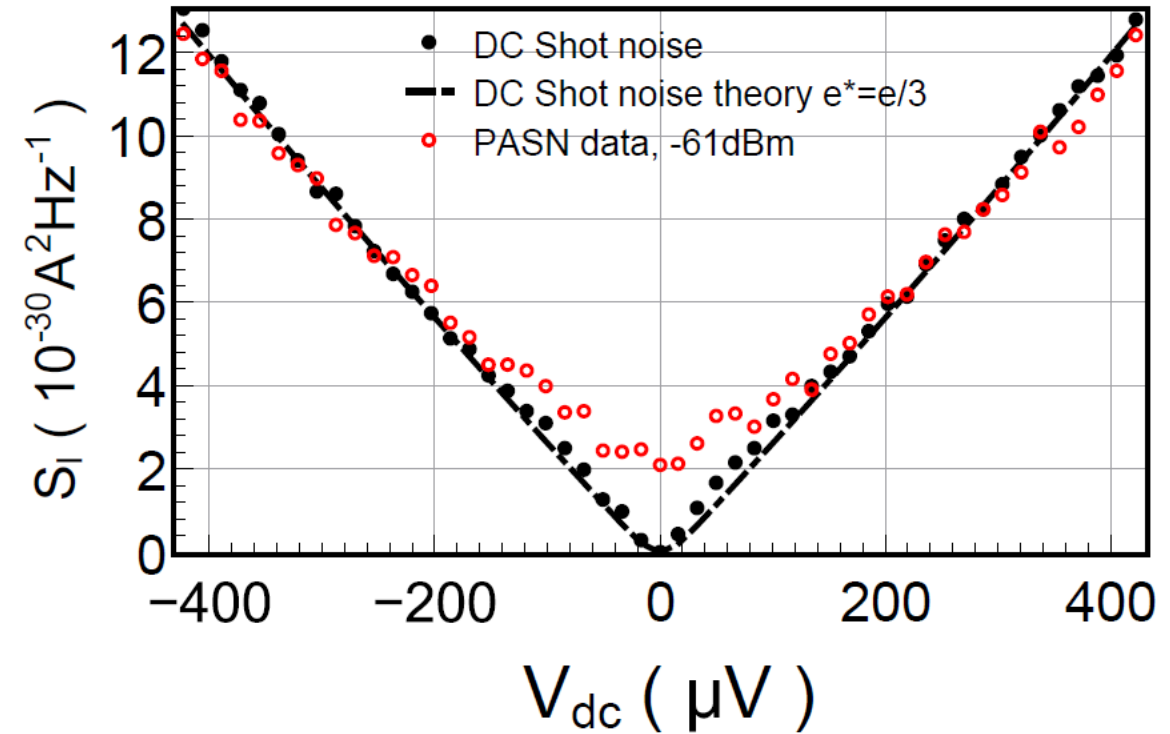
$$V(t) = V_{dc} + V_{ac} \cos(2\pi f t)$$

$$V_{ac} \approx 100 \mu\text{V} \text{ for } -67\text{dBm}$$

Photon-Assisted Shot Noise for the 1/3-FQHE state



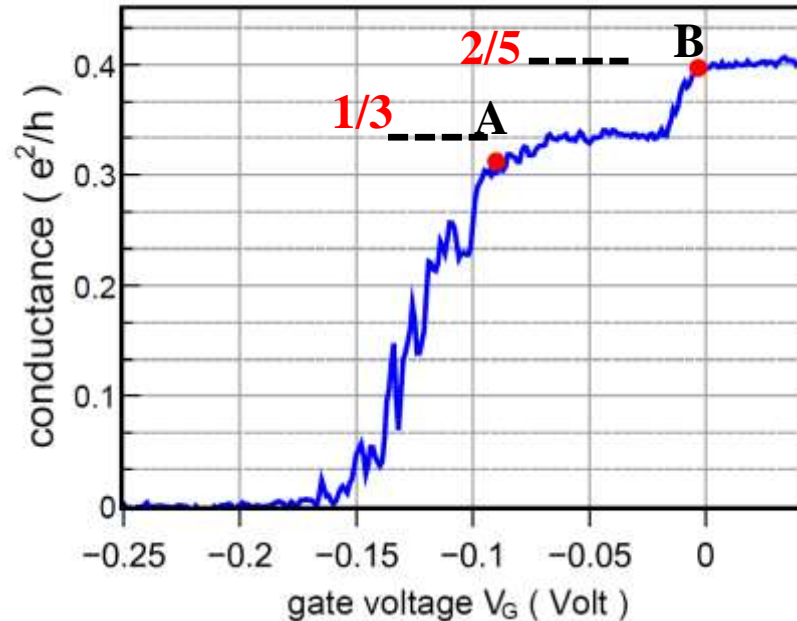
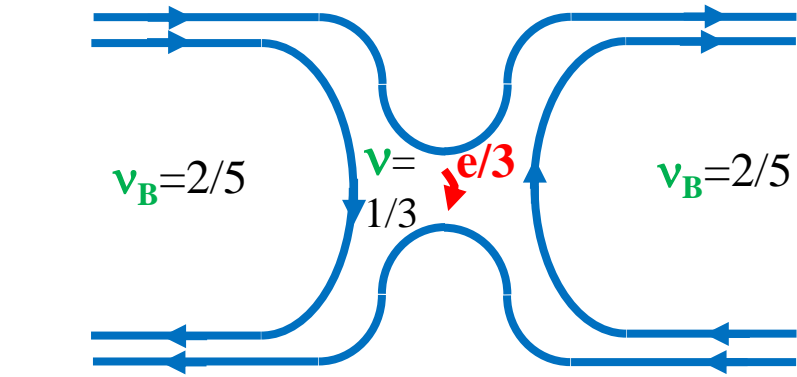
$f = 22 \text{ GHz}$



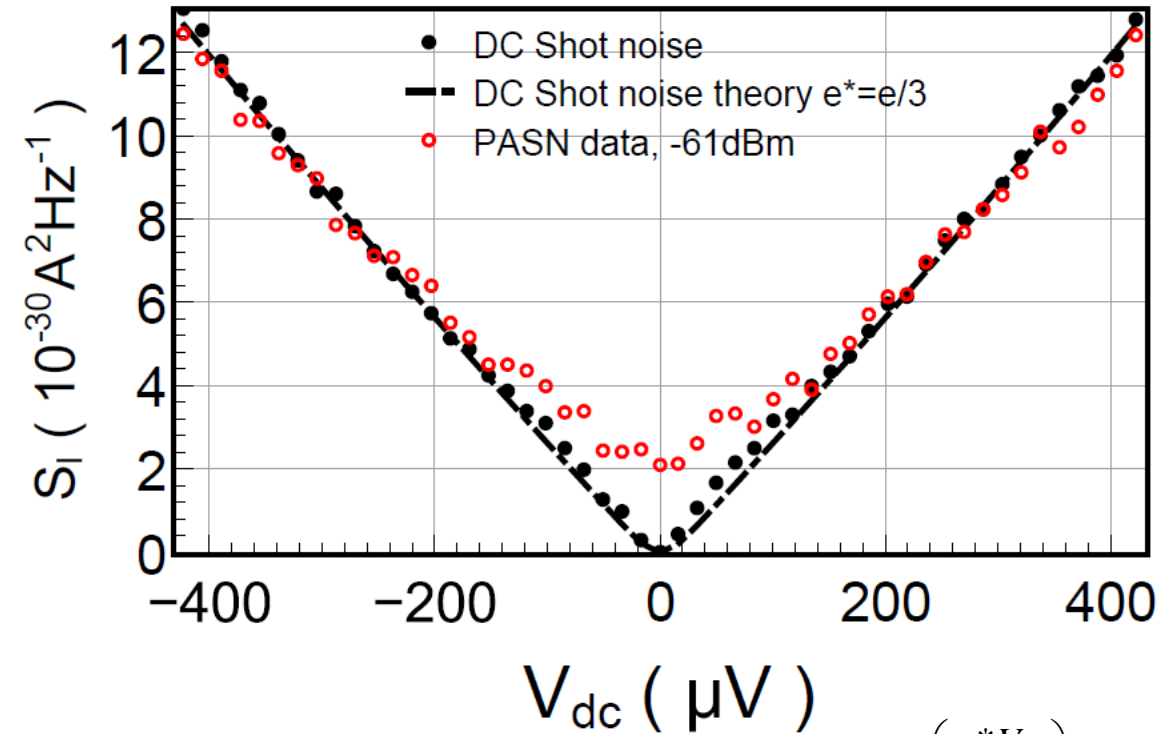
$$V(t) = V_{dc} + V_{ac} \cos(2\pi f t)$$

$$V_{ac} \approx 200 \mu\text{V} \text{ for } -61\text{dBm}$$

Photon-Assisted Shot Noise for the 1/3-FQHE state



$f=22\text{GHz}$



$$V(t) = V_{dc} + V_{ac} \cos(2\pi ft) \quad p_0 = J_0 \left(\frac{e^* V_{ac}}{hf} \right)$$

$$V_{ac} \approx 200 \mu\text{V} \text{ for } -61\text{dBm} \quad p_1 = -p_{-1} = J_1 \left(\frac{e^* V_{ac}}{hf} \right)$$

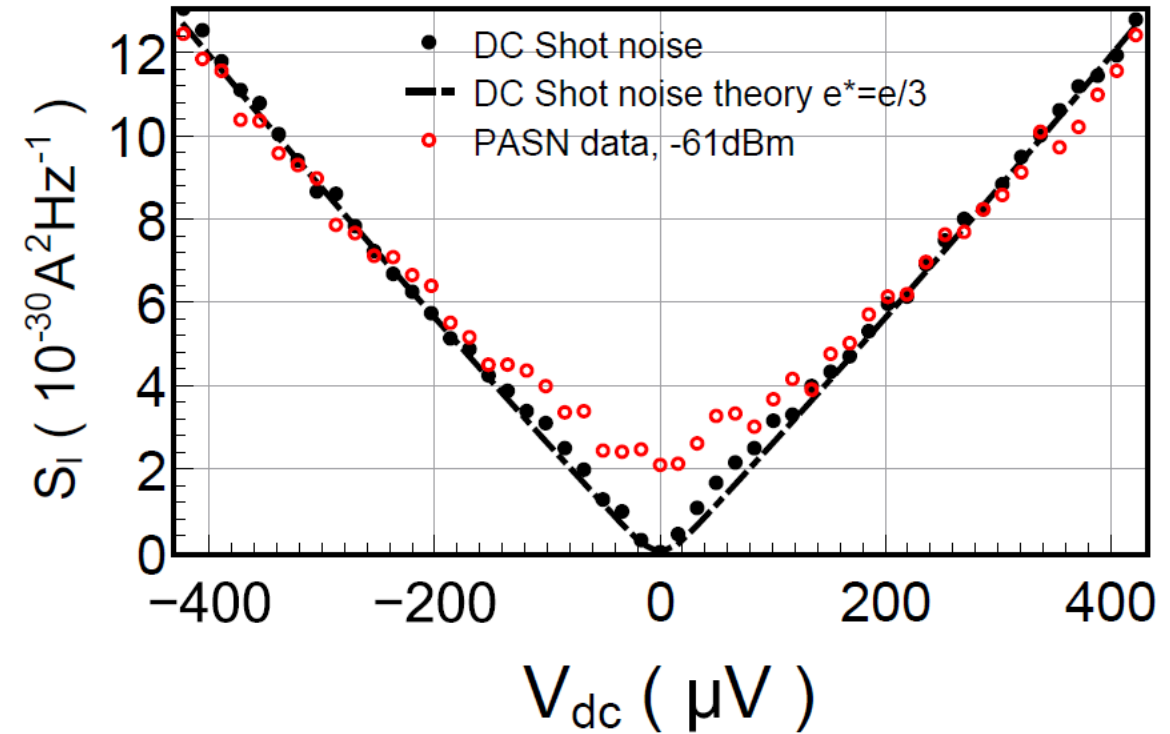
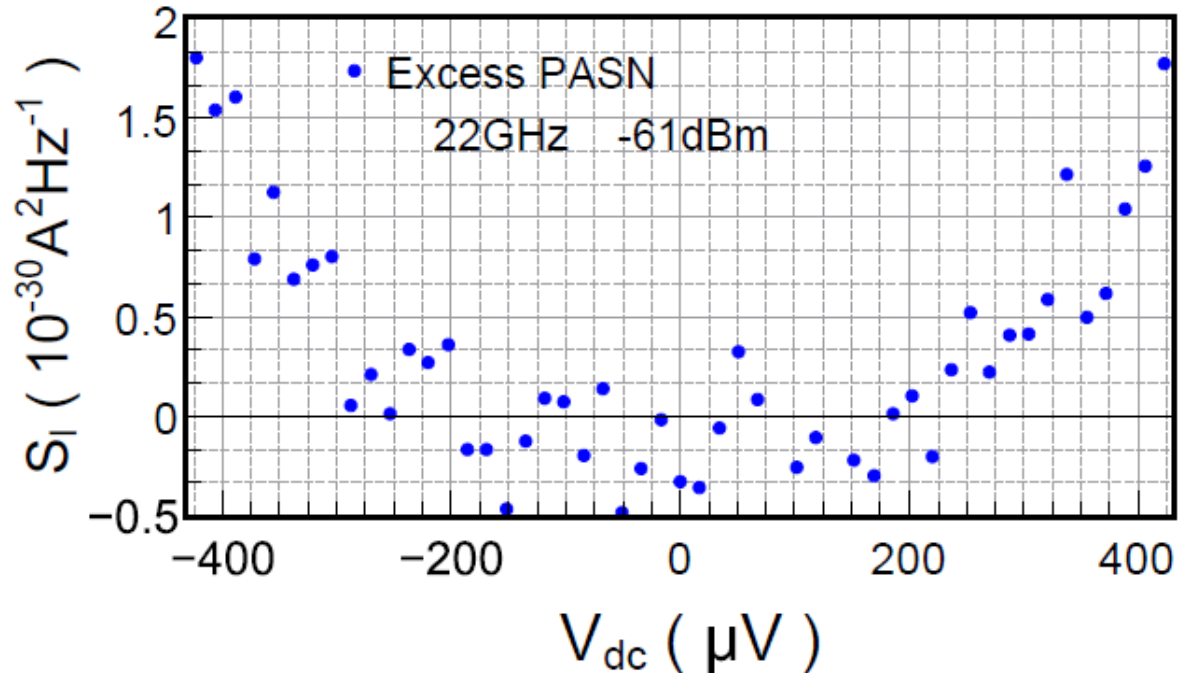
$$S_I^{PASN}(V_{dc}) \stackrel{?}{=} |p_0|^2 S_I^{DC}(V_{dc}) + |p_1|^2 \left[S_I^{DC}(V_{dc} - hf/e^*) + S_I^{DC}(V_{dc} + hf/e^*) \right]$$

Excess PASN for the 1/3-FQHE state

Killing the non photon-assisted part !

Excess PASN:

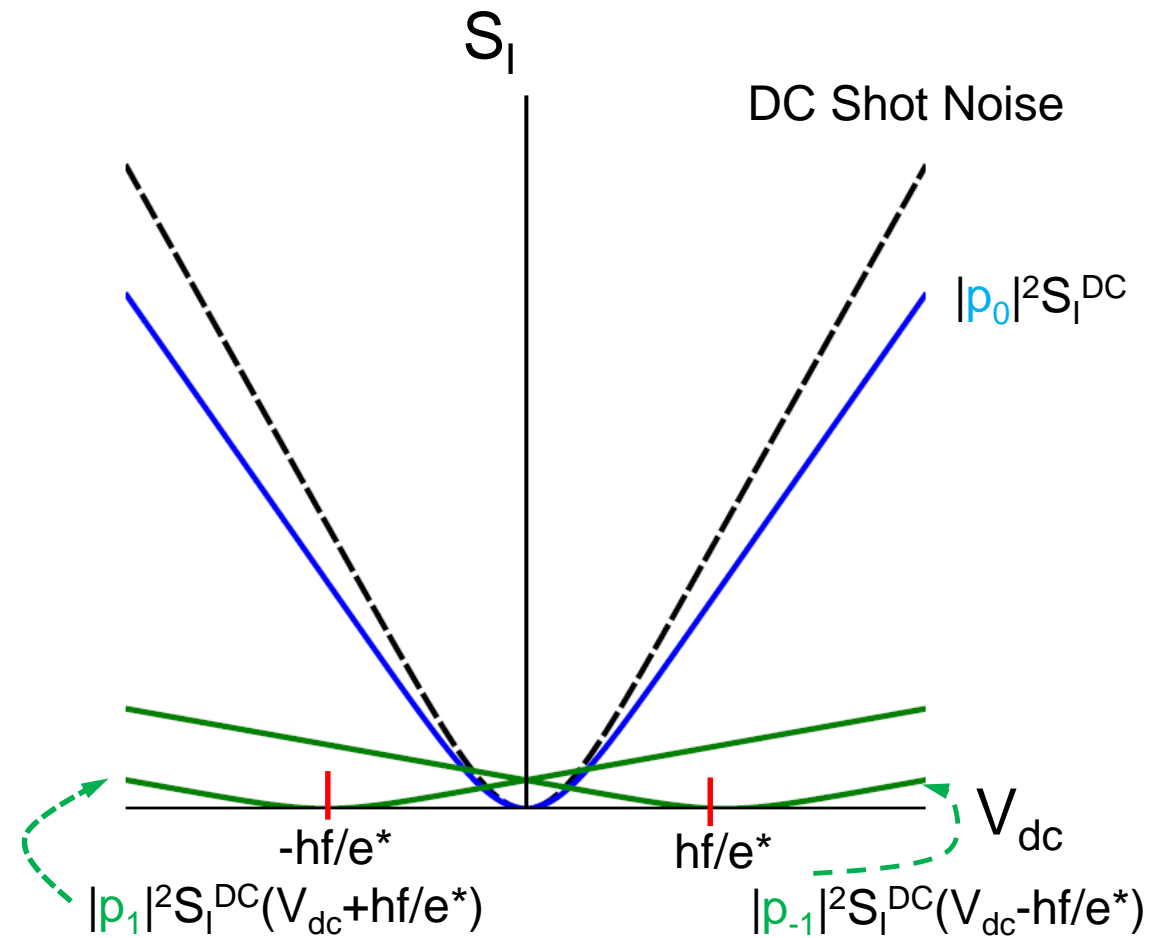
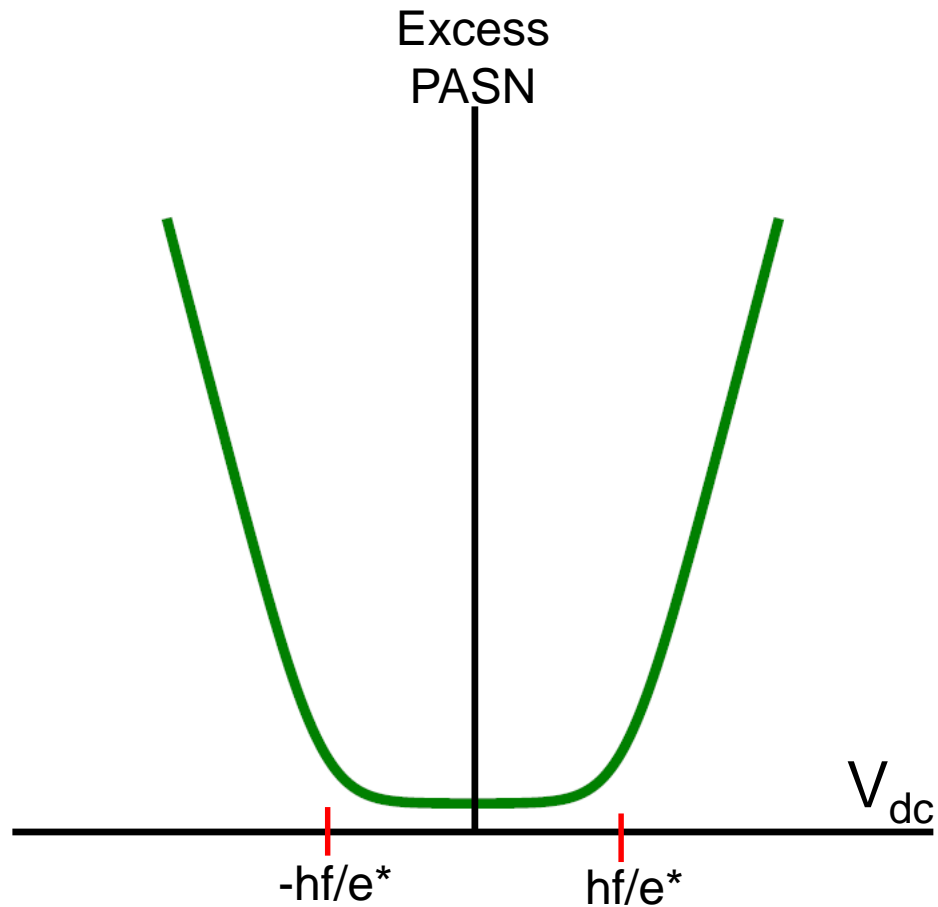
$$\begin{aligned}\Delta S_I &= S_I^{PASN}(V_{dc}) - |p_0|^2 S_I^{DC}(V_{dc}) \\ &= |p_1|^2 \left[S_I^{DC}(V_{dc} - hf/e^*) + S_I^{DC}(V_{dc} + hf/e^*) \right]\end{aligned}$$



Finding a *flat variation* for the low $|V_{dc}|$ range provides a determination of $|p_0|^2$

Excess PASN for the 1/3-FQHE state

WHY a FLAT VARIATION?

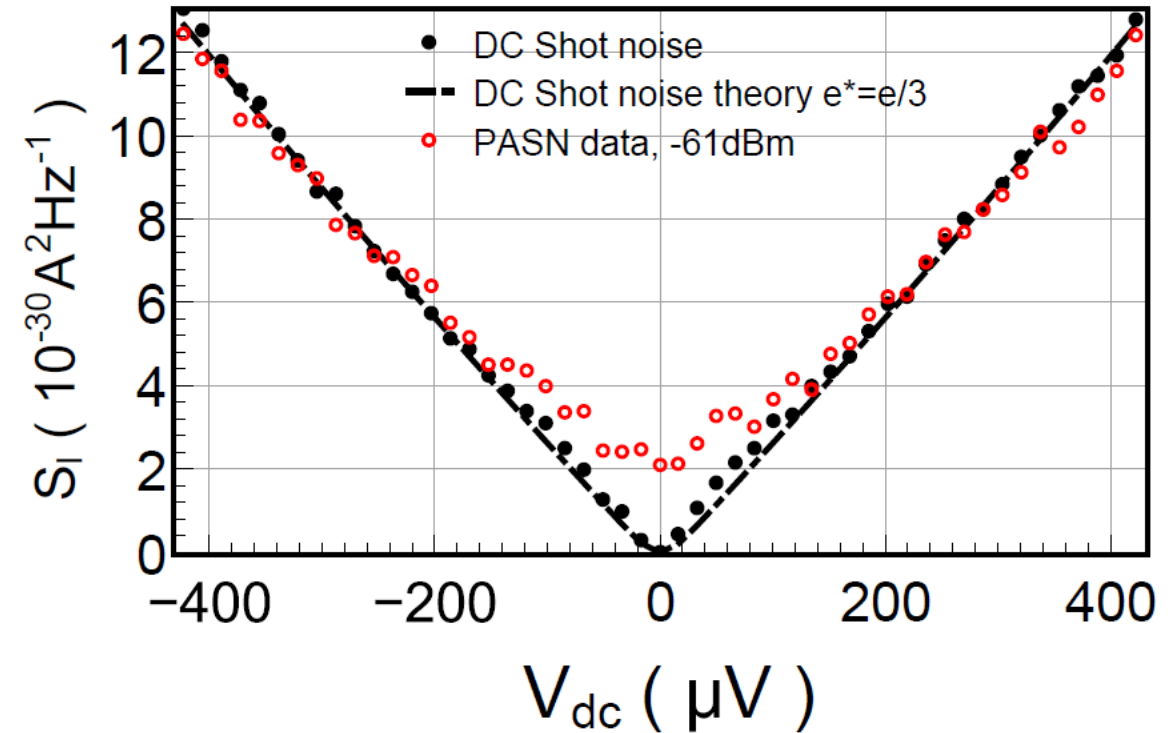
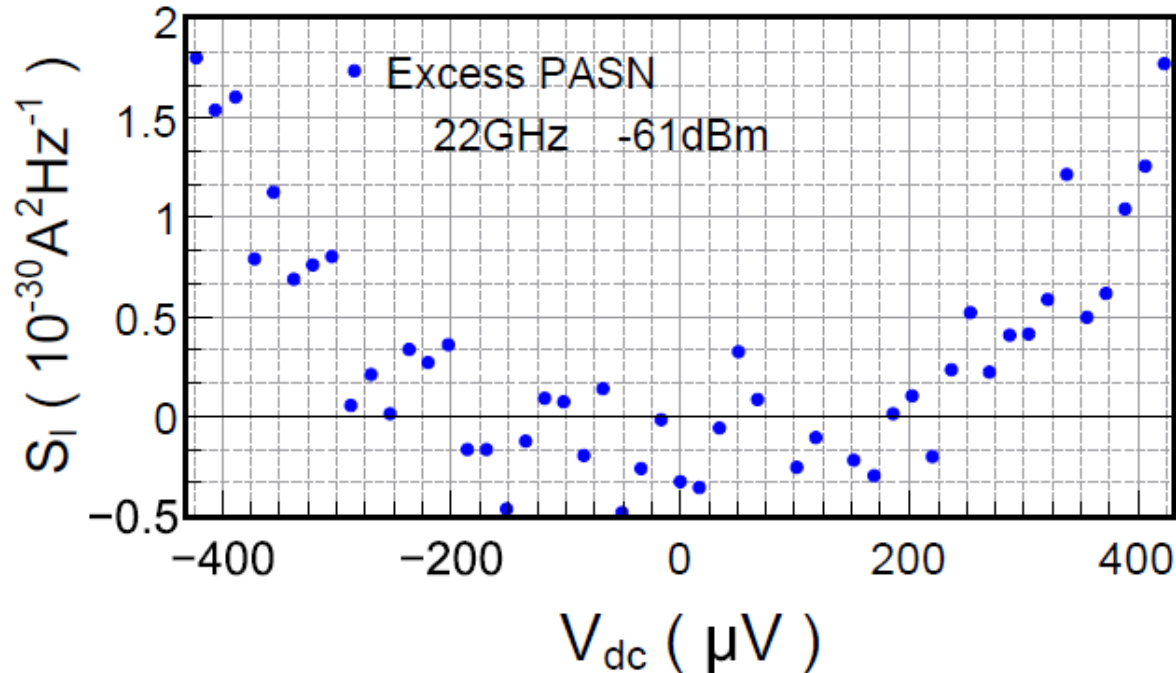


Excess PASN for the 1/3-FQHE state

Killing the non photon-assisted part !

Excess PASN:

$$\begin{aligned} \Delta S_I &= S_I^{PASN}(V_{dc}) - |p_0|^2 S_I^{DC}(V_{dc}) \\ &= |p_1|^2 \left[S_I^{DC}(V_{dc} - hf/e^*) + S_I^{DC}(V_{dc} + hf/e^*) \right] \end{aligned}$$



Finding a flat variation for the low $|V_{dc}|$ range provides a determination of $|p_0|^2$

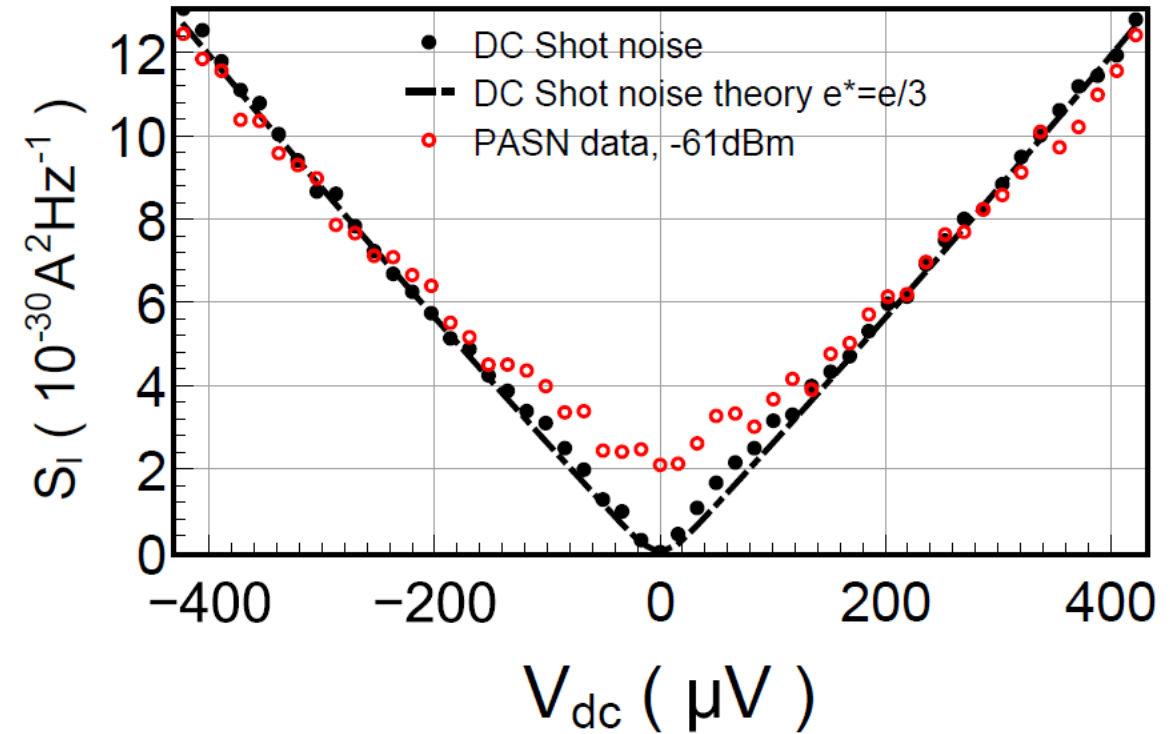
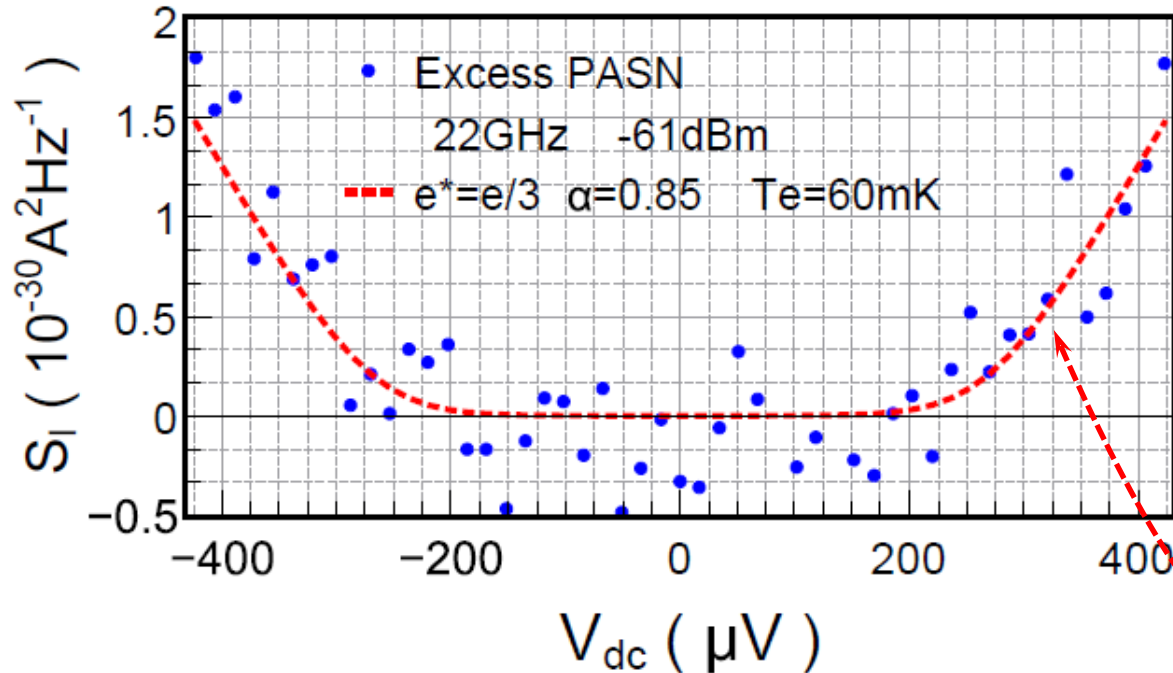
as $|p_0|^2 + 2|p_1|^2 \approx 1$ this gives $|p_1|^2$

Excess PASN for the 1/3-FQHE state

Killing the non photon-assisted part !

Excess PASN:

$$\begin{aligned} \Delta S_I &= S_I^{PASN}(V_{dc}) - |p_0|^2 S_I^{DC}(V_{dc}) \\ &= |p_1|^2 \left[S_I^{DC}(V_{dc} - hf/e^*) + S_I^{DC}(V_{dc} + hf/e^*) \right] \end{aligned}$$



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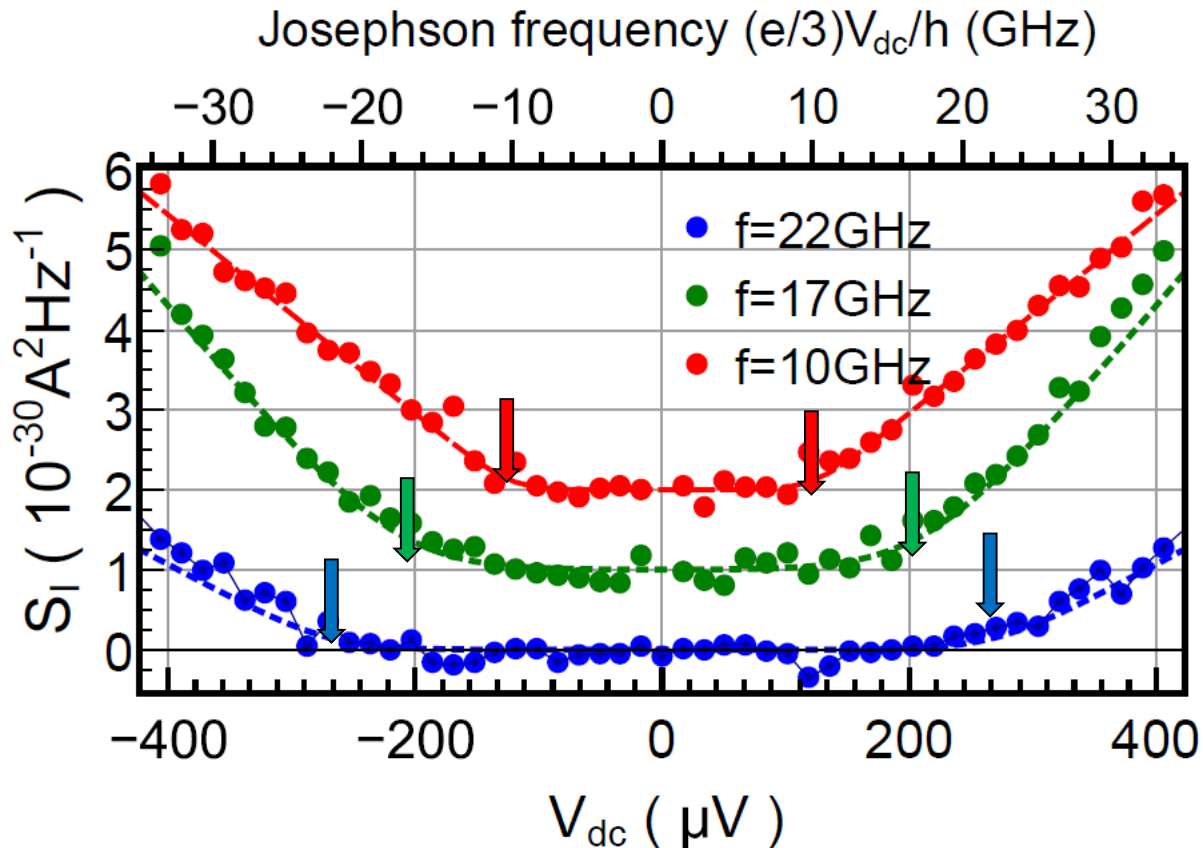
comparison using $f_{\text{Josephson}} = e^* V_{dc} / h$ with $e^* = e/3$

Josephson relation for the 1/3-FQHE state

CHECKING the FREQUENCY DEPENDENCE of Excess PASN:

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threshold voltage : $V_J = hf/e^*$ scales with frequency!



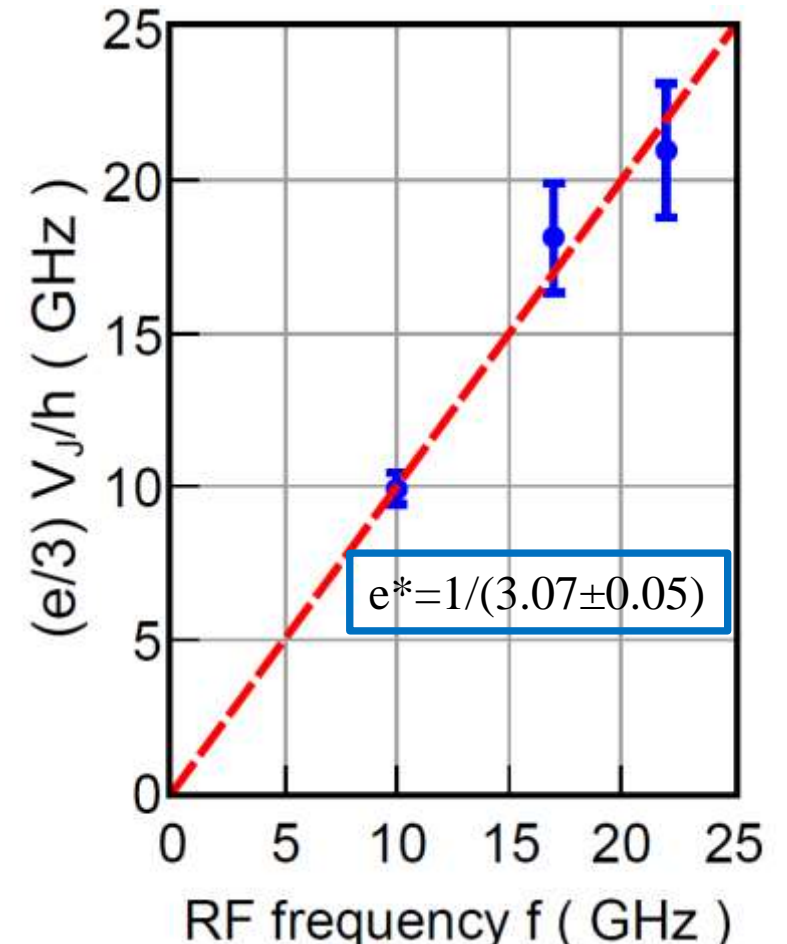
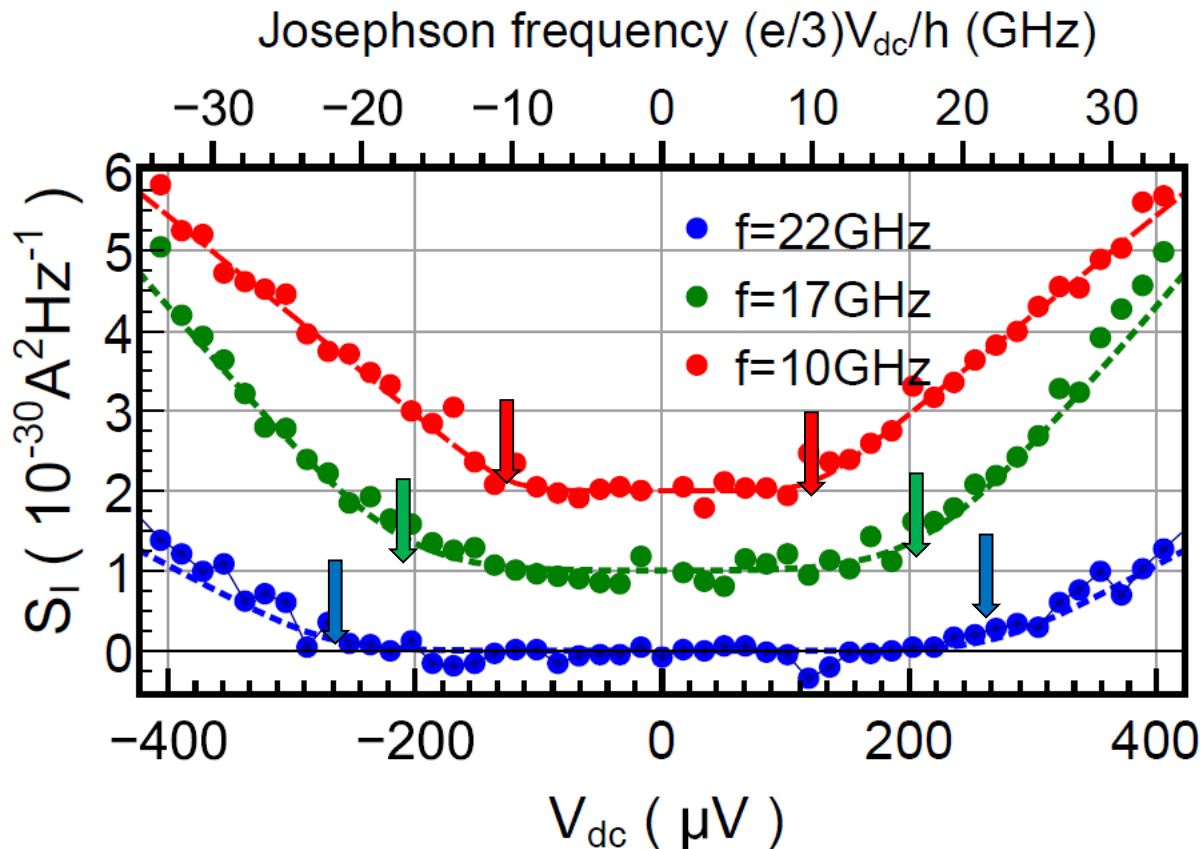
New Measurement of e^* for the 1/3-FQHE State

MEASURING e^* from Excess PASN:

$$\begin{aligned}\Delta S_I &= S_I^{PASN}(V_{dc}) - |p_0|^2 S_I^{DC}(V_{dc}) \\ &= |p_1|^2 \left[S_I^{DC}(V_{dc} - hf/e^*) + S_I^{DC}(V_{dc} + hf/e^*) \right]\end{aligned}$$

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Best fit of data with e^* free parameter



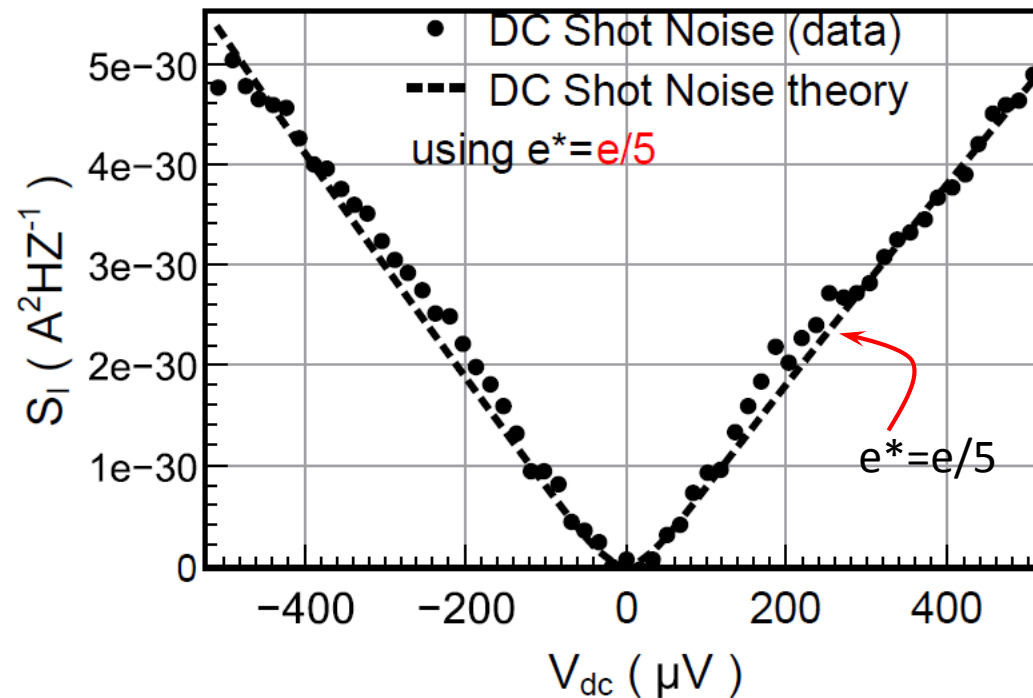
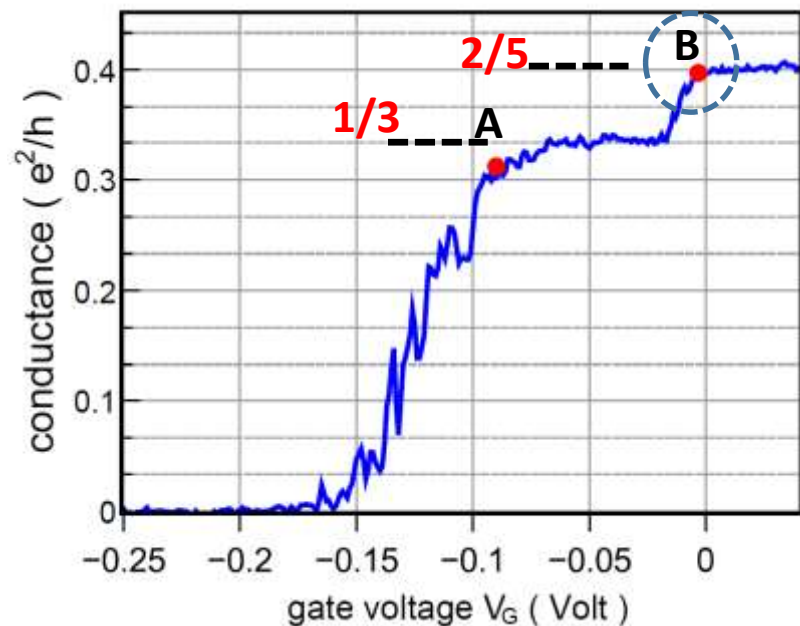
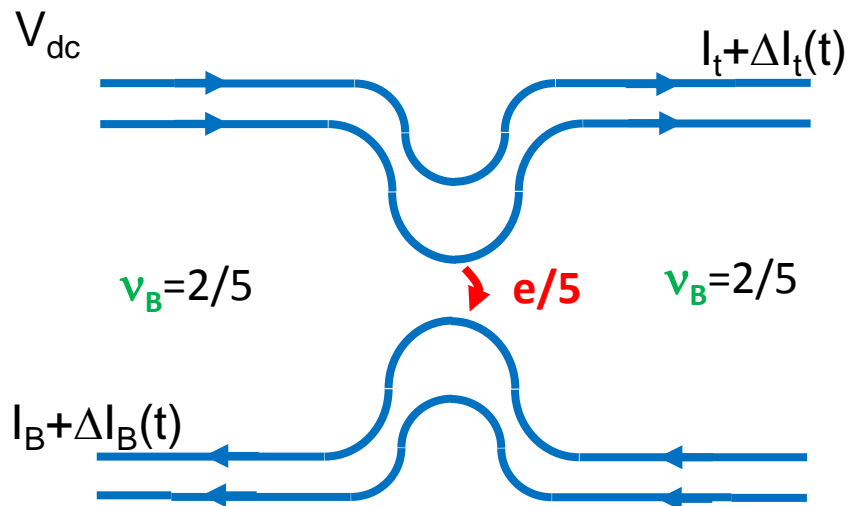
OUTLINE

- Quantum Hall edge states and Fractional Quantum Hall Effect
- PHOTON-ASSISTED TRANSPORT
 - Photon-assisted processes
 - A JOSEPHSON Relation for Photon Assisted Shot Noise (PASN)
- Experimental Results
 - $e^*=e/3$
 - $e^*=e/5$
- CONCLUSION and PERSPECTIVES

$$f_{J.} = \frac{e^*V}{h}$$

X. G. Wen (1991)

DC Shot noise for the 2/5-FQHE state

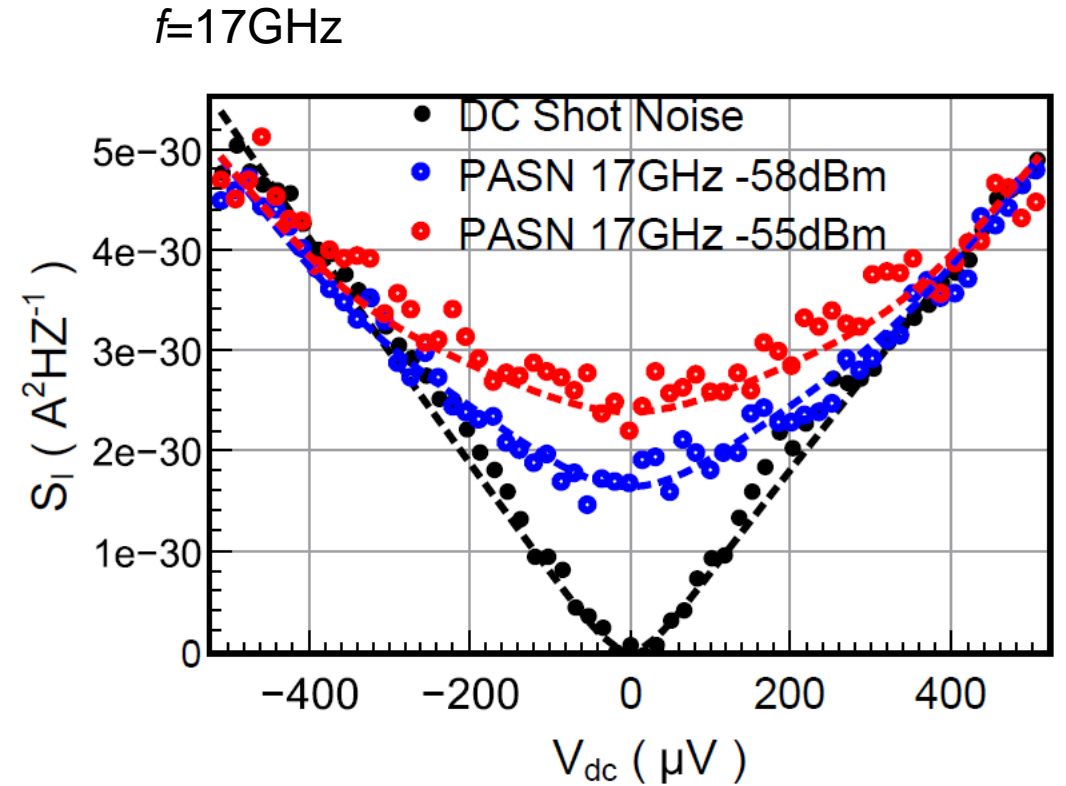
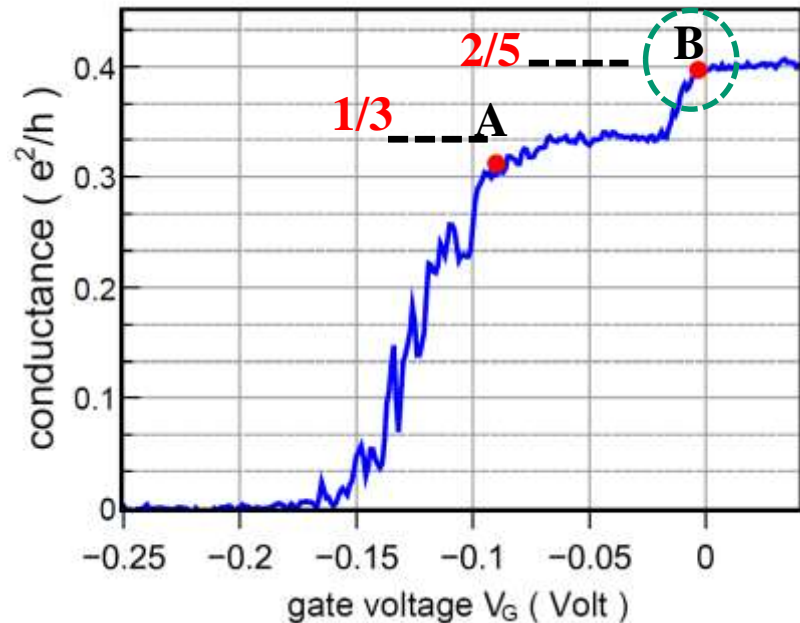
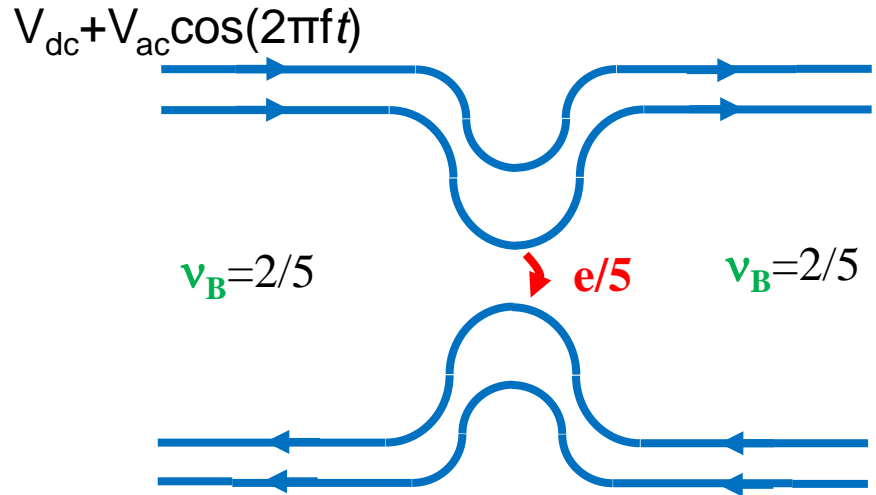


$$S_I^{DC} = 2e^* I_B \left[\coth \left(\frac{e^* V_{dc}}{2k_B T} \right) - \frac{2k_B T}{e^* V_{dc}} \right] \propto - \langle \Delta I_B \Delta I_t \rangle$$

$$e^* = e/5 !$$

confirms Weizmann results (Reznikov 1999) on 2/5

Photon-Assisted Shot Noise for the 2/5-FQHE state



$$V(t) = V_{dc} + V_{ac} \cos(2\pi ft)$$

$$V_{ac} \approx 300 \mu\text{V} \text{ for } -58\text{dBm}$$

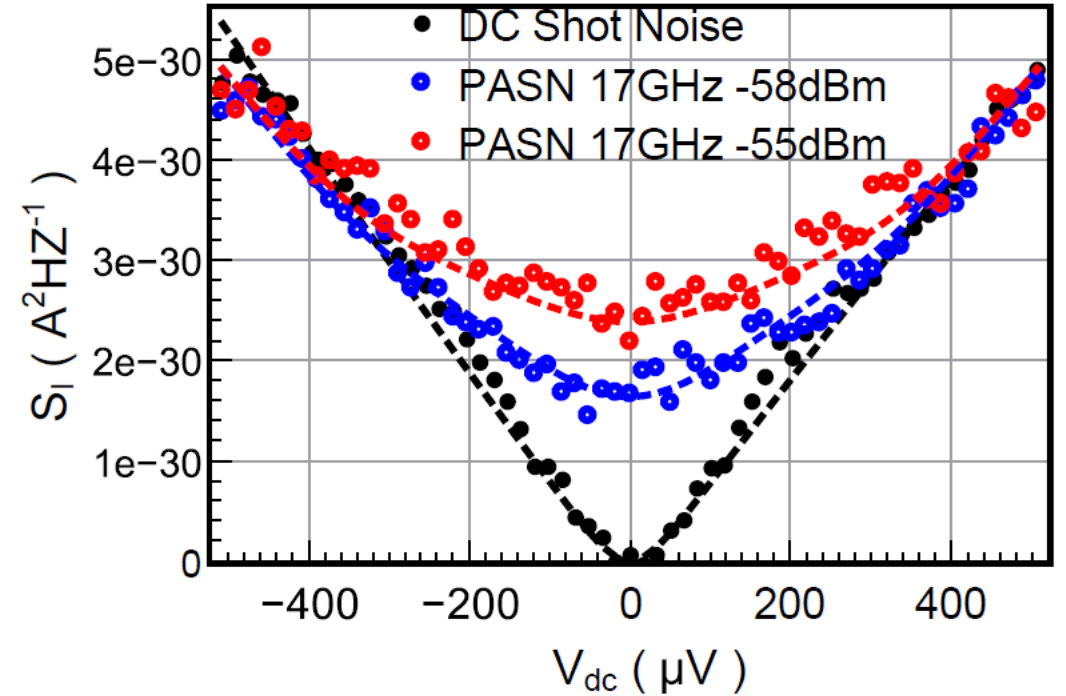
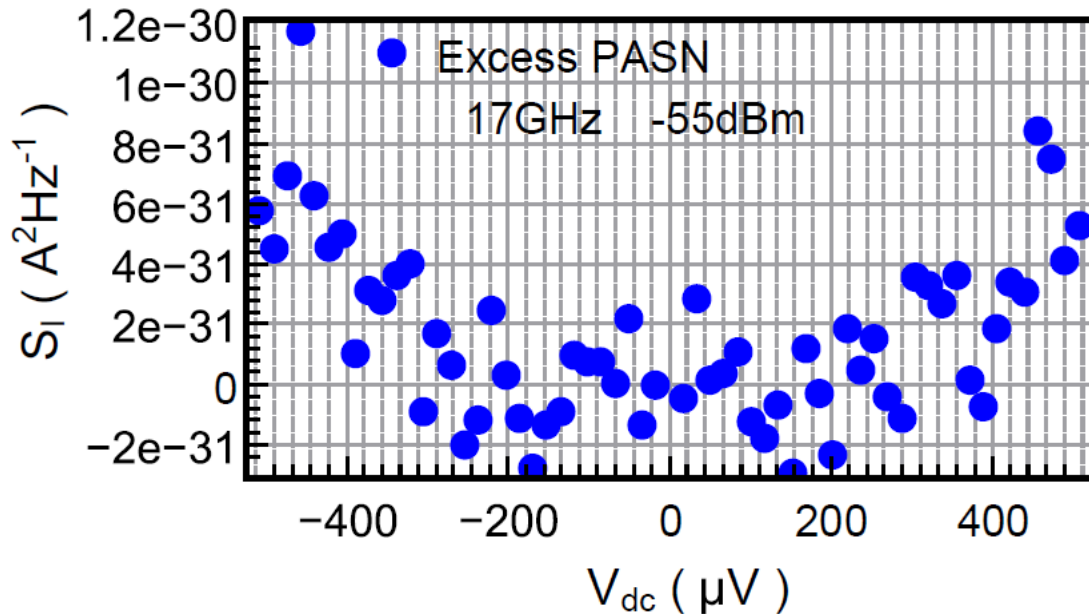
$$\approx 400 \mu\text{V} \text{ for } -55\text{dBm}$$

Excess PASN for the 2/5-FQHE state

Killing again the non photon-assisted part !

Excess PASN:

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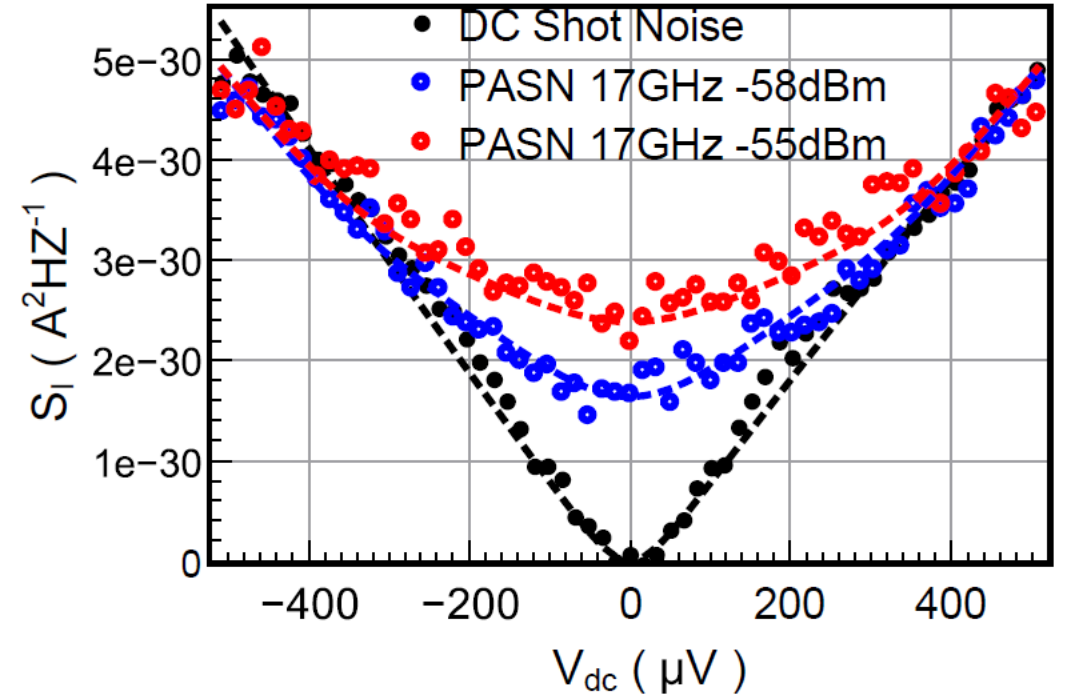
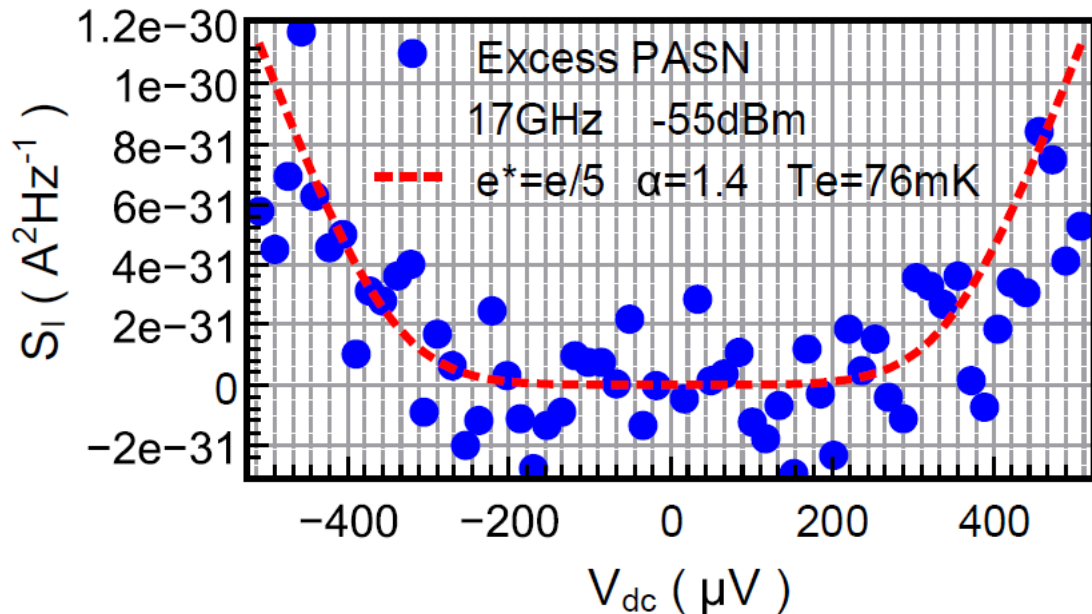
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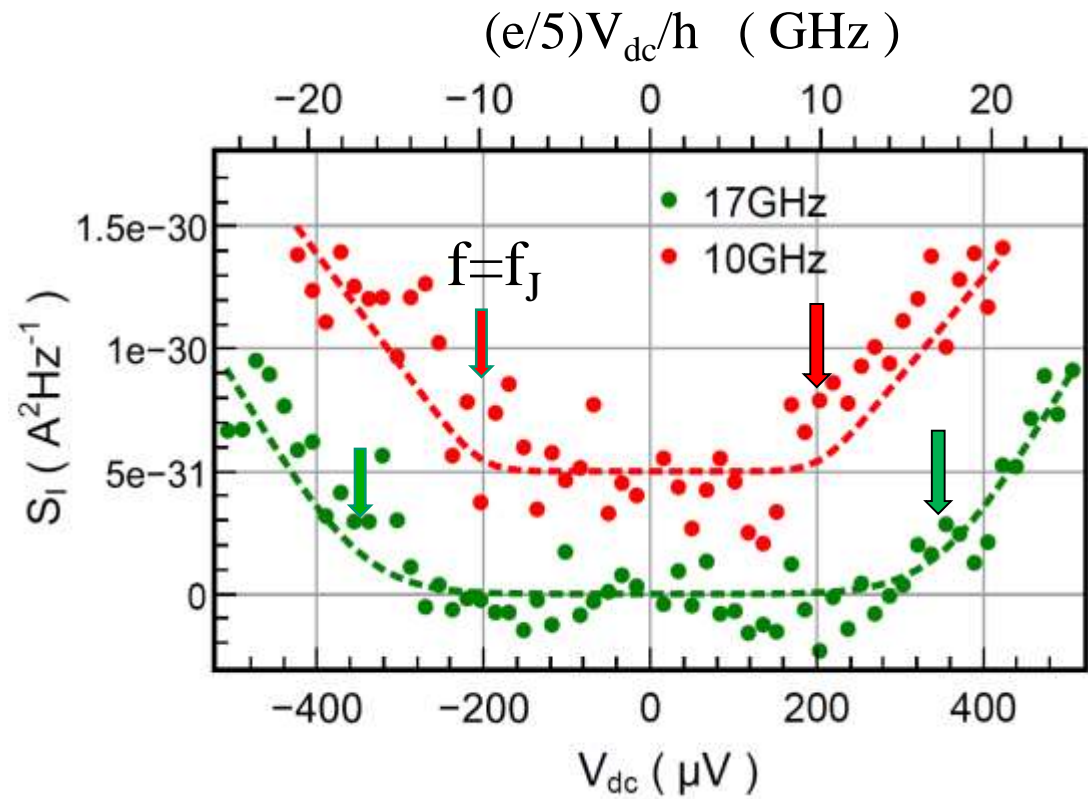
comparison using $f_{\text{Josephson}} = e^* V_{dc} / h$ with $e^* = e/5$

Josephson relation for the 2/5-FQHE state

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$$\begin{aligned}\Delta S_I &= S_I^{PASN}(V_{dc}) - |p_0|^2 S_I^{DC}(V_{dc}) \\ &= |p_1|^2 \left[S_I^{DC}(V_{dc} - hf/e^*) + S_I^{DC}(V_{dc} + hf/e^*) \right]\end{aligned}$$

threshold voltage : $V_J = hf/e^*$ scales with frequency!



New Measurement of e^* for the 2/5-FQHE State

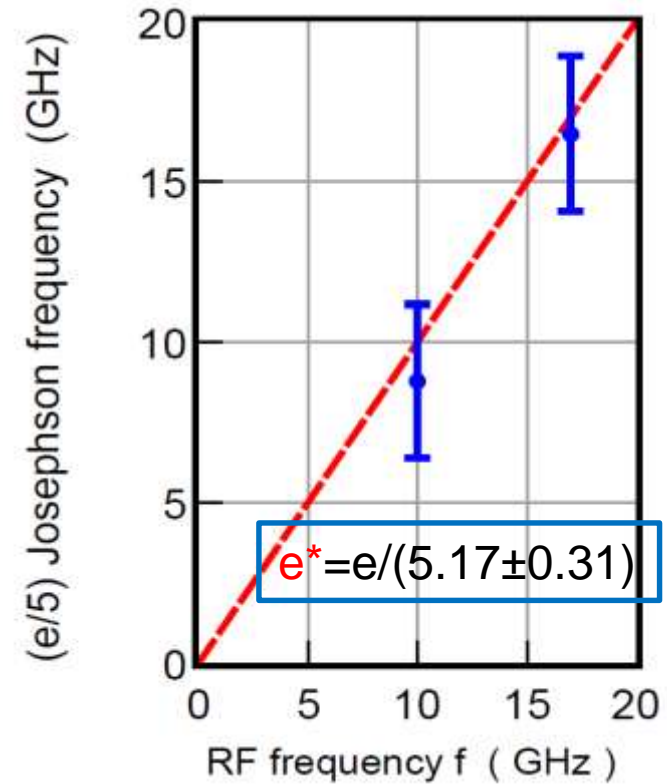
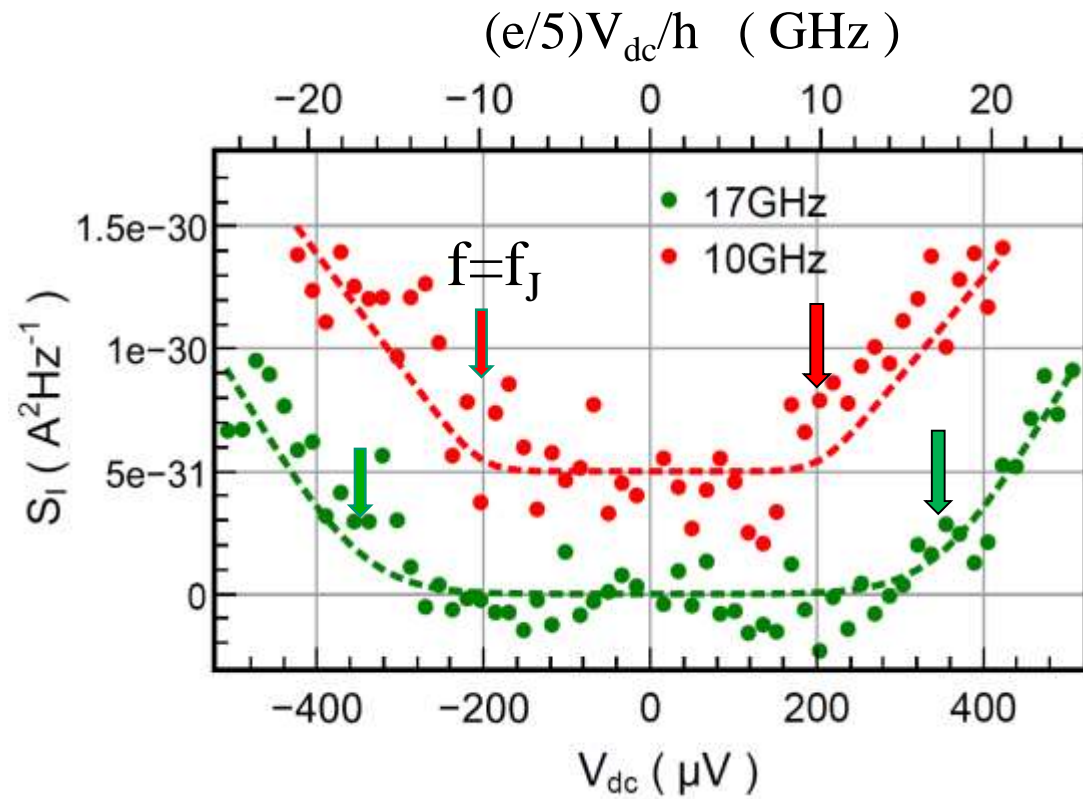
MEASURING e^* from Excess PASN:

$$\Delta S_I = S_I^{PASN}(V_{dc}) - |p_0|^2 S_I^{DC}(V_{dc})$$

$$= |p_1|^2 \left[S_I^{DC}(V_{dc} - hf/e^*) + S_I^{DC}(V_{dc} + hf/e^*) \right]$$

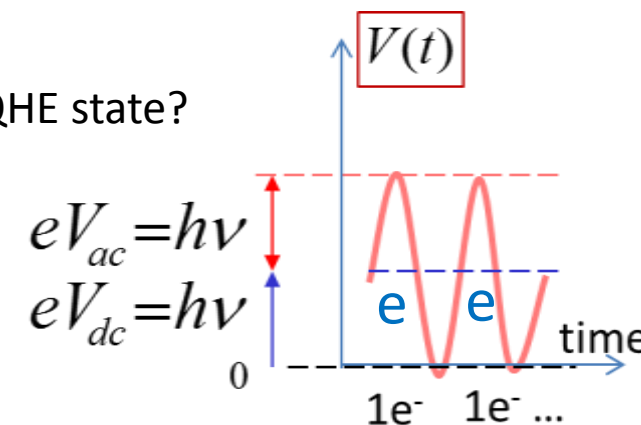
threshold voltage : $V_J = hf/e^*$ scales with frequency!

Best fit of data with e^* free parameter



CONCLUSION

- FQHE $e^*=e/3$ and $e/5$ abelian *anyons* can be manipulated with microwave by well-defined photon-assisted processes. What about $e/4$ in non-abelian $5/2$ FQHE state?
- Validates the possibility to realize on-demand single anyon sources for time domain *anyon braiding*.
- Based on Photon-Assisted Shot Noise (PASN)
- Shows evidence of the Josephson relation $e^*V/h=f$ predicted in 1991 by X.G. Wen*



Sine wave
single charge
pulses

(Old 1997 exp.)

SCHOTTKY (charge granularity)

$$S_I = 2 e^* I_B$$

weak signal
but accurate

good signal
but lack of accuracy,
model dependent

(New 2019 exp.)

PASN Josephson Relation (photon quantum)

$$h f = e^* V$$

very accurate

good accuracy

*predicted for the current, see also I. Safi +Sukhorukov (2010).

ACKNOWLEDGEMENTS



OPEN POSITION
for 18-24 months
Post-doct.
(urgent)

X. Waintal

H. Saleur

I. Safi

Th. Martin

M. Freedman

All members of Nanoelectronics Group at Saclay

The cryogeny Team

ANR FullyQuantum AAP CE30

The Josephson Frequency of fractionally charge anyons
M. Kapfer, P. Roulleau, I. Farrer, D. A. Ritchie, and D. C. Glattli,
[arXiv:1806.03117](https://arxiv.org/abs/1806.03117),

Published 24 January 2019 on **Science**

DOI: [10.1126/science.aau3539](https://doi.org/10.1126/science.aau3539)

Levitons :

J. Dubois et al, Nature 502, 659 (2013)

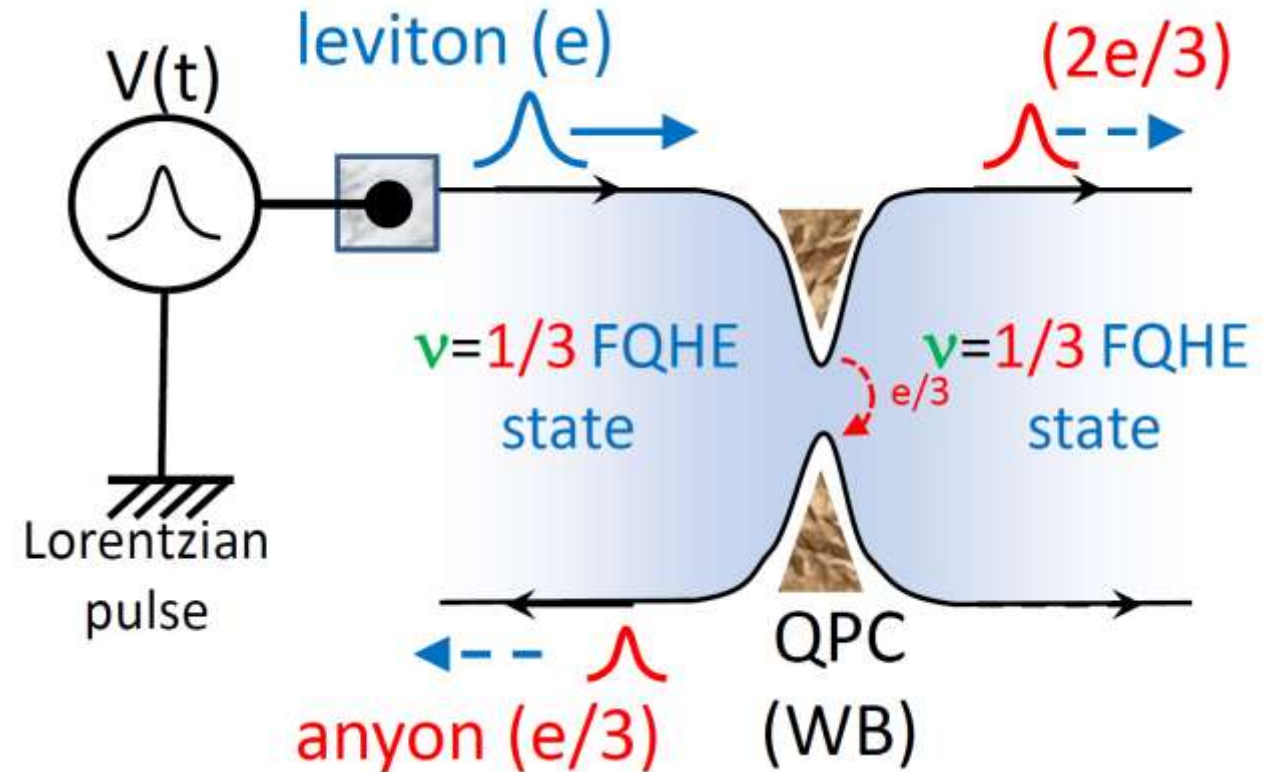
T. Jullien et al., Nature 514, 603 (2014)

PERSPECTIVE : ANYONS on DEMAND

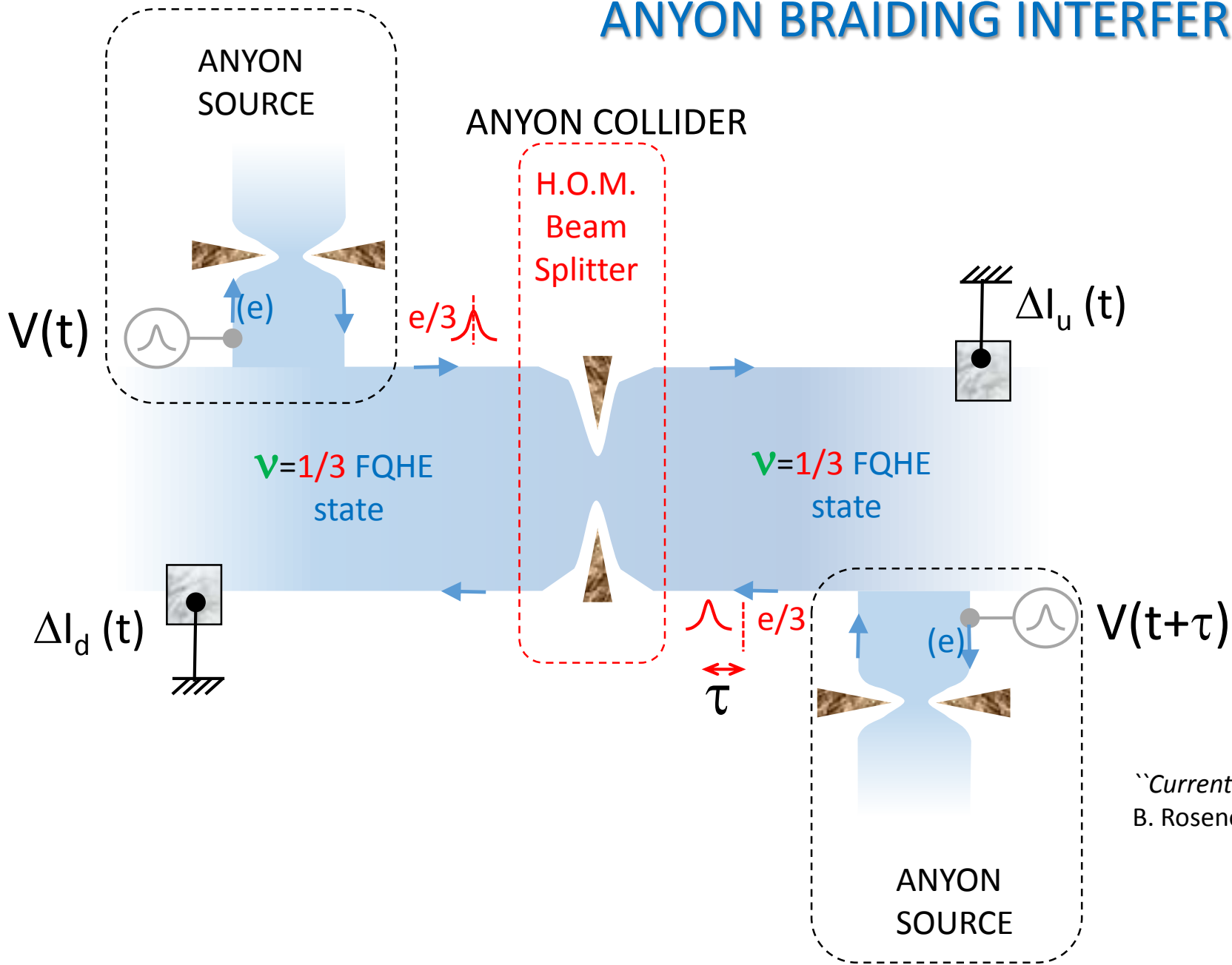
A Time Controlled Poissonian Source of Anyon

IDEA: Weak backscattering beaks the leviton into $e/3$, $2e/3$ quasiparticles.

- Anyons inherit from the time properties of Levitons
- Non-deterministic: Poissonian source



PERSPECTIVE : ANYON BRAIDING INTERFERENCE

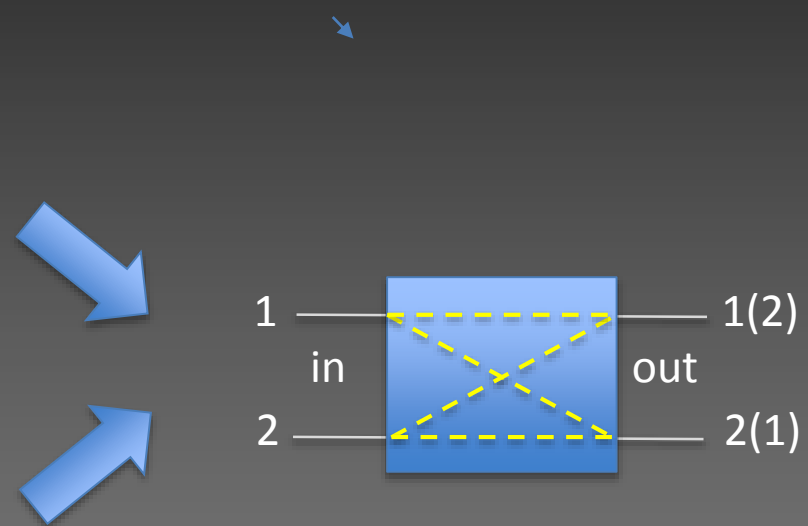
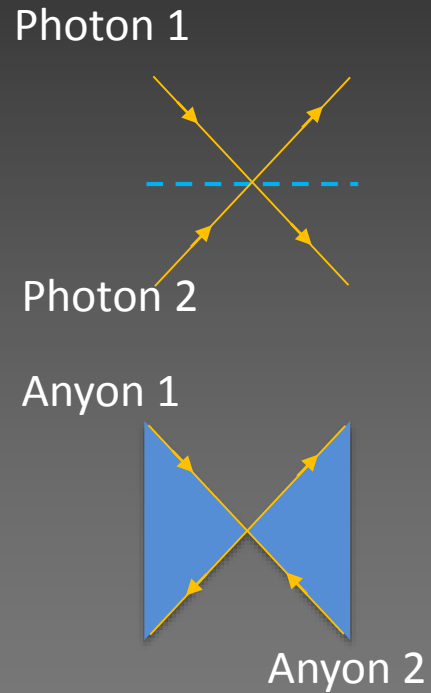


$$-\langle \Delta I_u \Delta I_d \rangle \propto 1 + g_2(\tau) \cos(\theta_{stat.})$$

“Current Correlations from a Mesoscopic Anyon Collider”
B. Rosenow, I. P. Levkivskyi, B. I. Halperin, (2016)

Braiding Anyons

1) Unveiling the anyon statistical angle with Hong Ou Mandel braiding interference



$$P(1,2) = |b(1,2)+b(2,1)|^2$$

$$P(1,2) = (1 - \cos\theta)/2$$

- 0 : boson bunching ($\theta=0$)
- 1 : fermion antibunching ($\theta=\pi$)
- $1/4$: for $\nu=1/3$ FQHE abelian anyons ($\theta=\pi/3$)